

EQUILIBRIUM STATES IN DYNAMICAL SYSTEMS VIA GEOMETRIC MEASURE THEORY

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ABSTRACT. Given a dynamical system with a uniformly hyperbolic (chaotic) attractor, the physically relevant Sinai–Ruelle–Bowen (SRB) measure can be obtained as the limit of the dynamical evolution of the leaf volume along local unstable manifolds. We extend this geometric construction to the substantially broader class of equilibrium states corresponding to Hölder continuous potentials; these states arise naturally in statistical physics and play a crucial role in studying stochastic behavior of dynamical systems. The key step in our construction is to replace leaf volume with a reference measure that is obtained from a Carathéodory dimension structure via an analogue of the construction of Hausdorff measure. In particular, we give a new proof of existence and uniqueness of equilibrium states that does not use standard techniques based on Markov partitions or the specification property; our approach can be applied to systems that do not have Markov partitions and do not satisfy the specification property.

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1. INTRODUCTION

1.1. Systems with hyperbolic behavior. A *smooth dynamical system with discrete time* consists of a smooth manifold M —the *phase space*—and a diffeomorphism $f: M \rightarrow M$. Each state of the system is represented by a point $x \in M$, whose *orbit* $(f^n(x))_{n \in \mathbb{Z}}$ gives the time evolution of that state. We are interested in the case where the dynamics of f exhibit *hyperbolic* behavior. Roughly speaking,

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this means that orbits of nearby points separate exponentially quickly in either forward or backward time; if the phase space is compact, this leads to the phenomenon popularly known as *chaos*.

Hyperbolic behavior turns out to be quite common, and for such systems it is not feasible to make specific forecasts of a single trajectory far into the future, because small initial errors quickly grow large enough to spoil the prediction. On the other hand, one may hope to make statistical predictions about the asymptotic behavior of orbits of f . A measurement of the system corresponds to a function $\varphi: M \rightarrow \mathbb{R}$; the sequence $\varphi, \varphi \circ f, \varphi \circ f^2, \dots$ represents the same observation made at successive times. When specific forecasts of $\varphi \circ f^n$ are impossible, we can treat this sequence as a stochastic process and make predictions about its asymptotic behavior. For a more complete discussion of this point of view, see [ER85], [Mañ87, Chapter 1], and [Via97].

1.2. Physical measures and equilibrium states. To fully describe the stochastic process $(\varphi \circ f^n)_{n \in \mathbb{Z}}$, we need a probability measure μ on M that represents the likelihood of finding the system in a given state at the present time. The measure $f_*\mu$ defined by $\int \varphi d(f_*\mu) = \int \varphi \circ f d\mu$ represents the distribution one unit of time into the future. An *invariant* measure has $\mu = f_*\mu$, and hence $\mu = f_*^n\mu$ for all n , so the sequence of observations becomes a *stationary* stochastic process.

In this paper we will consider *uniformly hyperbolic systems*, for which the tangent bundle admits an invariant splitting $TM = E^u \oplus E^s$ such that E^u is uniformly expanded and E^s uniformly contracted by Df ; see §2 for examples and §3.1 for a precise definition. Such systems have an extremely large set of invariant measures; for example, standard results show that there are infinitely many periodic orbits, each supporting an atomic invariant measure. Thus one is led to the problem of selecting a distinguished measure, or class of measures, that is most dynamically significant.

Since we work on a smooth manifold, it would be natural to consider an invariant volume form on M , or at least an invariant measure that is absolutely continuous with respect to volume. However, for dissipative systems such as the *solenoid* described in §2.2, no such invariant measure exists, and one must instead look for a *Sinai–Ruelle–Bowen (SRB) measure*, which we describe in §§3.2–3.3. Such a measure is absolutely continuous “in the unstable direction”, which is enough to guarantee that it is *physically relevant*; it describes the asymptotic statistical behavior of volume-typical trajectories.

SRB measures can be constructed via the following *geometric approach*: let m be a normalized Lebesgue measure (volume) for some Riemannian metric on M , consider its forward iterates $f_*^n m$, then average the first N of these and take a limit measure as $N \rightarrow \infty$.

Another approach to SRB measures, which we recall in §3.4, is via *thermodynamic formalism*, which imports mathematical tools from equilibrium statistical physics in order to describe the behavior of large ensembles of trajectories. This program began in the late 1950s, when Kolmogorov and Sinai introduced the concept of entropy into dynamical systems; see [Kat07] for a historical overview. Given a *potential function* $\varphi: M \rightarrow \mathbb{R}$, one studies the *equilibrium states* associated to φ , which are invariant measures that maximize the quantity $h_\mu(f) + \int \varphi d\mu$, where h_μ

denotes the Kolmogorov–Sinai entropy.¹ The maximum value is called the *topological pressure* of φ and is denoted $P(\varphi)$.

In the 1960s and 1970s, it was shown by Sinai, Ruelle, and Bowen that for uniformly hyperbolic systems, every Hölder continuous potential has a unique equilibrium state (see §3.1). Applying this result to the particular case of the *geometric potential*² $\varphi^{\text{geo}}(x) = -\log |\det Df|_{E^u(x)}|$, one has $P(\varphi^{\text{geo}}) = 0$ and the equilibrium state is the SRB measure described above; see §3.4.

1.3. Different approaches to constructing equilibrium states. There are two main classical approaches to thermodynamic formalism. The first uses *Markov partitions* of the manifold M ; we recall the general idea in §3.5. The second approach uses the *specification property*, which we review in §3.6.

The purpose of this paper is to describe a third *geometric* approach, which was outlined above for SRB measures: produce an equilibrium state as a limiting measure of the averaged pushforwards of some *reference measure*, which need not be invariant. For the physical SRB measure, this reference measure was Lebesgue; to extend this approach to other equilibrium states, one must start by choosing a new reference measure. The definition of this reference measure, and its motivation and consequences, is the primary goal of this paper, and our main result can be roughly stated as follows.

For every Hölder continuous potential φ , one can use the tools of geometric measure theory to define a reference measure m_φ for which the averages of the pushforwards $f_^n m_\varphi$ converge to the unique equilibrium state for φ .*

A precise statement of the result is given in §4. An important motivation for this work is that the geometric approach can be applied to more general situations beyond the uniformly hyperbolic systems studied in this paper. For example, the geometric approach was used in [CDP16] to construct SRB measures for some nonuniformly hyperbolic systems, and in [CPZ18] we use it to construct equilibrium states for some partially hyperbolic systems. The first two approaches—Markov partitions and specification—have also been extended beyond uniform hyperbolicity (see [CP17] for a survey of the literature), but the overall theory in this generality is still very far from being complete, so it seems worthwhile to add another tool by developing the geometric approach as well.

1.4. Reference measures for general potentials. In the geometric construction of the physical SRB measure, one can take the reference measure to be either a Lebesgue measure m on M or a Lebesgue measure m_W on any *local unstable leaf* $W = V_{\text{loc}}^u(x)$. These leaves are d^u -dimensional submanifolds of M that are tangent at each point to the *unstable distribution* $E^u(x) \subset T_x M$; they are expanded by the dynamics of f and have the property that $f(V_{\text{loc}}^u(x)) \supset V_{\text{loc}}^u(f(x))$; see §3.1 for more details. Given a local unstable leaf $W = V_{\text{loc}}^u(x)$, we will write m_W or m_x^u for the

¹From the statistical physics point of view, the quantity $E_\mu := -(h_\mu(f) + \int \varphi d\mu)$ is the free energy of the system, so that an equilibrium state minimizes the free energy; see [Sar15, §1.6] for more details.

²Here the determinant is taken with respect to any orthonormal bases for $E^u(x)$ and $E^u(f(x))$. If the map f is of the class of smoothness $C^{1+\alpha}$ for some $\alpha > 0$, then one can show that $\varphi^{\text{geo}}(x)$ is Hölder continuous.

leaf volume determined by the induced Riemannian metric. This has the following key properties.

- (1) $m_W = m_x^u$ is a finite nonzero Borel measure on W .
- (2) If W_1 and W_2 are local unstable leaves with nontrivial intersection, then m_{W_1} and m_{W_2} agree on the overlap.
- (3) Under the dynamics of f , the leaf volumes scale by the rule

$$(1.1) \quad m_{f(x)}^u(A) = \int_{f^{-1}A} |\det Df|_{E^u(y)}| dm_x^u(y).$$

As mentioned in §1.2, the geometric potential $\varphi^{\text{geo}}(x) = -\log |\det Df|_{E^u(x)}|$ has $P(\varphi^{\text{geo}}) = 0$, so the integrand in (1.1) can be written as $e^{P(\varphi^{\text{geo}}) - \varphi^{\text{geo}}(y)}$. In §4.3, given a continuous potential φ , we will construct on every local unstable leaf $W = V_{\text{loc}}^u(x)$ a reference measure $m_x^{\mathcal{C}}$ satisfying similar properties to m_x^u , but with the scaling rule³

$$(1.2) \quad m_{f(x)}^{\mathcal{C}}(A) = \int_{f^{-1}A} e^{P(\varphi) - \varphi(y)} dm_x^{\mathcal{C}}(y).$$

The superscript \mathcal{C} is shorthand for a *Carathéodory dimension structure* determined by the potential φ and a scale $r > 0$; see §5 for the essential facts about such structures, and see [Pes97] for a complete description. Roughly speaking, the definitions of $P(\varphi)$ and $m_x^{\mathcal{C}}$ are analogous to the definitions of Hausdorff dimension and Hausdorff measure, respectively, but they take the dynamics into account. Recall that the latter definitions involve covers by balls of decreasing radius; the modification to obtain our quantities involves covering by dynamically defined balls, as explained in §4.3.

1.5. Some history. The idea of constructing dynamically significant measures for uniformly hyperbolic maps by first finding measures on unstable leaves with certain scaling properties goes at least as far back as work of Sinai [Sin68], which relies on Markov partitions. For uniformly hyperbolic dynamical systems with continuous time (flows) and the potential $\varphi = 0$, the corresponding equilibrium state, which is the measure of maximal entropy, was obtained by Margulis [Mar70]. He used a different construction of leaf measures via functional analysis of a special operator (induced by the dynamics) acting on the Banach space of continuous functions with compact support on unstable leaves. These leaf measures were studied further in [RS75, BM77].

Hasselblatt gave a description of the Margulis measure in terms of Hausdorff dimension [Has89], generalizing a result obtained by Hamenstädt for geodesic flows on negatively curved compact manifolds [Ham89]. In this geometric setting, where stable and unstable leaves are naturally identified with the ideal boundary of the universal cover, Kaimanovich observed in [Kai90, Kai91] that these leaf measures could be identified with the measures on the ideal boundary introduced by Patterson [Pat76] and Sullivan [Sul79]. For geodesic flows in negative curvature, this approach was recently extended to nonzero potentials by Paulin, Pollicott, and Schapira [PPS15].

For general hyperbolic systems and nonzero potential functions, families of leaf measures with the appropriate scaling properties were constructed by Haydn [Hay94] and Leplaideur [Lep00], both using Markov partitions. The key innovation in the

³Note that in general $P(\varphi) \neq 0$.

present paper is that we can construct these leaf measures directly, without using Markov partitions, by an approach reminiscent of Hasselblatt's from [Has89]. This requires us to interpret quantities in thermodynamic formalism by analogy with Hausdorff dimension, an idea which was introduced by Bowen for entropy in [Bow73], developed by Pesin and Pitskel' for pressure in [PP84], and generalized further by Pesin in [Pes88, Pes97].

1.6. Plan of the paper. We describe some motivating examples in §2, and give general background definitions in §3. These sections are addressed to a general mathematical audience, and the reader who is already familiar with thermodynamic formalism for hyperbolic dynamical systems can safely skip to §4, where we give the new definition of the reference measures m_x^C and formulate our main results. In §5 we recall the necessary results on Carathéodory dimension characteristics and describe some applications of our results to dimension theory. For well-known general results, we omit the proofs and give references to the literature where proofs can be found. For the new results stated here, we give an outline of the proofs in §6 and refer to [CPZ18] for complete details.

2. MOTIVATING EXAMPLES

Before recalling general definitions about uniformly hyperbolic systems and their invariant measures in §3, we describe three examples to motivate the idea of a *physical measure*. Our discussion here is meant to convey the overall picture and omits many details.

2.1. Hyperbolic toral automorphisms. Our first example is the diffeomorphism f on the torus $\mathbb{T}^2 = \mathbb{R}^2/\mathbb{Z}^2$ induced by the linear action of the matrix $L = \begin{pmatrix} 2 & 1 \\ 1 & 1 \end{pmatrix}$ on \mathbb{R}^2 , as shown in Figure 2.1.

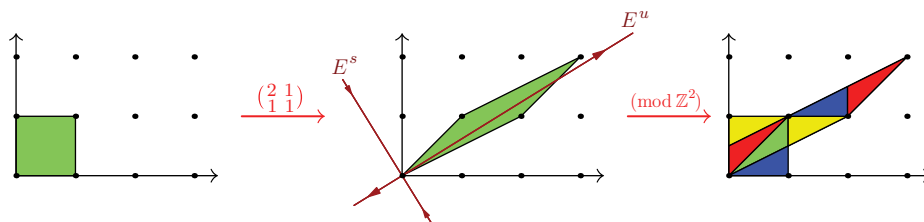


FIGURE 2.1. Lebesgue measure is preserved by f

This system is *uniformly hyperbolic*: The matrix L has two positive real eigenvalues $\lambda_s < 1 < \lambda_u$, whose associated eigenspaces E^s and E^u give a Df -invariant splitting of the tangent bundle $T\mathbb{T}^2$. The lines in \mathbb{R}^2 parallel to these subspaces project to f -invariant foliations W^s and W^u of the torus.

What about invariant measures? If $p \in \mathbb{T}^2$ has $f^n(p) = p$, then the measure $\mu = \frac{1}{n}(\delta_p + \delta_{f(p)} + \cdots + \delta_{f^{n-1}(p)})$ is invariant. Every point with rational coordinates is periodic for f , so this gives infinitely many f -invariant measures. Lebesgue measure is also invariant since $\det Df = \det L = 1$. (This is far from a complete list, as we will see.)

A measure μ is *ergodic* if every f -invariant function (every $\varphi \in L^1(\mu)$ with $\varphi = \varphi \circ f$) is constant μ -almost everywhere. One can check easily that the periodic

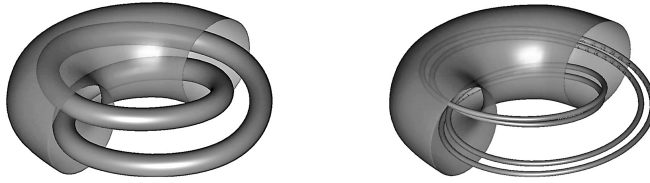


FIGURE 2.2. No absolutely continuous invariant measure

orbit measures from above are ergodic and, with a little more work, that Lebesgue measure is ergodic too.⁴ Birkhoff’s ergodic theorem says that if μ is ergodic, then μ -almost everywhere orbit has asymptotic behavior controlled by μ . More precisely, we say that the *basin of attraction* for μ is the set of initial conditions satisfying a law of large numbers governed by μ for continuous observables:

$$(2.1) \quad B_\mu = \left\{ x \in \mathbb{T}^2 : \frac{1}{n} \sum_{k=0}^{n-1} \varphi(f^k x) \xrightarrow{n \rightarrow \infty} \int \varphi d\mu \text{ for all } \varphi \in C(\mathbb{T}^2, \mathbb{R}) \right\}.$$

The ergodic theorem says that if μ is ergodic, then $\mu(B_\mu) = 1$.

For periodic orbit measures, this says very little, since it leaves open the possibility that the measure μ only controls the asymptotic behavior of finitely many orbits.⁵ For Lebesgue measure m , however, this says quite a lot: m governs the statistical behavior of Lebesgue-almost every orbit, and in particular, a point chosen at random with respect to any volume form on \mathbb{T}^2 has a trajectory whose asymptotic behavior is controlled by m . This is the sense in which Lebesgue measure is the *physically relevant* invariant measure, and we make the following definition.

Definition 2.1. An invariant measure μ for a diffeomorphism f is a *physical measure* if its basin B_μ has positive volume.

2.2. Smale–Williams solenoid. From Birkhoff’s ergodic theorem, we see that if μ is an ergodic invariant measure that is equivalent to a volume form,⁶ then that volume form gives full weight to the basin B_μ , and so a volume-typical trajectory has asymptotic behavior controlled by μ .

The problem now is that there are many examples for which no such μ exists. One such is the *Smale–Williams solenoid* studied in [Sma67, §I.9] and [Wil67]; see also [PC09, Lecture 29] for a gentle introduction and further discussion. This is a map from the open solid torus U into itself. Abstractly, the solid torus is the direct product of a disc and a circle, so that one may use coordinates (x, y, θ) on U , where x and y give coordinates on the disc and θ is the angular coordinate on the circle. Define a map $f : U \rightarrow U$ by

$$(2.2) \quad f(x, y, \theta) := \left(\frac{1}{4}x + \frac{1}{2} \cos \theta, \frac{1}{4}y + \frac{1}{2} \sin \theta, 2\theta \right).$$

Figure 2.2 shows two iterates of f , with half of the original torus for reference.

Every invariant measure is supported on the *attractor* $\Lambda = \bigcap_{n \geq 0} f^n(U)$, which has zero volume. In particular, there is no invariant measure that is absolutely

⁴This can be proved either by Fourier analysis or by the more geometric *Hopf argument*; see §4.2.2.

⁵In fact B_μ is infinite, being a union of leaves of the stable foliation W^s .

⁶Recall that two measures μ and ν are *equivalent* if $\mu \ll \nu$ and $\nu \ll \mu$, in which case we write $\mu \sim \nu$.

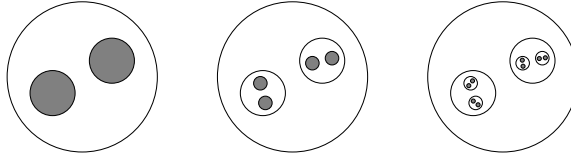


FIGURE 2.3. A cross-section of the attractor

continuous with respect to volume. Nevertheless, it is still possible to find an invariant measure that is physically relevant in the sense given above. To do this, first observe that since the solenoid map f contracts distances along each cross-section $\mathbb{D}^2 \times \{\theta_0\}$, any two points in the same cross-section have orbits with the same (forward) asymptotic behavior: given an invariant measure μ , the basin B_μ is a union of such cross-sections.

This fact suggests that we should look for an invariant measure that is absolutely continuous in the direction of the circle coordinate θ , which is expanded by f . To construct such a measure, observe that each cross-section intersects the images $f^n(U)$ in a nested sequence of unions of discs (as shown in Figure 2.3) so the attractor Λ intersects this cross-section in a Cantor set. Thus Λ is locally the direct product of an interval in the expanding direction and a Cantor set in the contracting directions. Let m^u be the Lebesgue measure on the circle, and let μ be the measure on Λ that projects to m^u and gives equal weight to each of the 2^n pieces at the n th level of the Cantor set construction in Figure 2.3. One can show without too much difficulty that μ is invariant and ergodic and that, moreover, B_μ has full volume in the solid torus U . Thus even though μ is singular, it is still the physically relevant invariant measure due to its absolute continuity in the expanding direction.

2.3. Smale’s horseshoe. Finally, we recall an example for which no physical measure exists—the *horseshoe* introduced by Smale in the early 1960s; see [Sma67, §I.5] and [PC09, Lecture 31] for more details, see also [Sma98] for more history. Consider a map $f: R \rightarrow \mathbb{R}^2$ which acts on the square $R := [0, 1]^2$ as shown in Figure 2.4: first the square is contracted vertically by a factor of $\alpha < 1/2$ and stretched horizontally by a factor of $\beta > 2$, then it is bent and positioned so that $f(R) \cap R$ consists of two rectangles of height α and length 1.

Observe that a part of the square R is mapped to the complement of R . Consequently, f^2 is not defined on the whole square R but only on the union of two vertical strips in R . The set where f^3 is defined is the union of four vertical strips, two inside each of the previous ones, and so on; there is a Cantor set $C^u \subset [0, 1]$

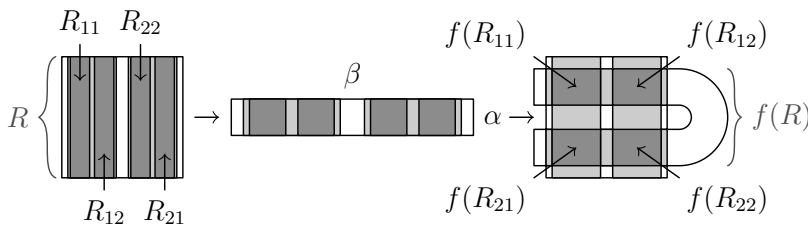


FIGURE 2.4. No physical measure

such that every point outside $C^u \times [0, 1]$ can be iterated only finitely many times before leaving R . In particular, every f -invariant measure has $B_\mu \subset C^u \times [0, 1]$ and, hence, B_μ is Lebesgue-null, so there is no physical measure.

Note that the argument in the previous paragraph did not consider the stable (vertical) direction at all. For completeness, observe that there is a Cantor set $C^s \subset [0, 1]$ such that $\bigcap_{n \geq 0} f^n(R) = [0, 1] \times C^s$ and that the maximal f -invariant set $\Lambda := \bigcap_{n \in \mathbb{Z}} f^n(R)$ is a direct product $C^u \times C^s$.

2.4. Main ideas. The three examples discussed so far have certain features in common, which are representative of uniformly hyperbolic systems.

First, every invariant measure lives on a compact invariant set Λ that is locally the direct product of two sets, one contracted by the dynamics and one expanded. For the hyperbolic toral automorphism $\begin{pmatrix} 2 & 1 \\ 1 & 1 \end{pmatrix}$, $\Lambda = \mathbb{T}^2$, and both of these sets were intervals; for the solenoid, there was an interval in the expanding direction and a Cantor set in the contracting direction; for the horseshoe, both were Cantor sets.

Second, the physically relevant invariant measure (when it existed) could also be expressed as a direct product. For the hyperbolic toral automorphism, it was a product of Lebesgue measure on the two intervals. For the solenoid, it was a product of Lebesgue measure on the interval (the expanding circle coordinate) and a $(\frac{1}{2}, \frac{1}{2})$ -Bernoulli measure on the contracting Cantor set.

Third, and most crucially for our purposes, in identifying the physical measure, it is enough to look at how invariant measures behave *along the expanding (unstable) direction*. We will make this precise in §3.2 when we discuss *conditional measures*, and this idea will motivate our main construction in §4.3 of reference measures associated to different potential functions.

Note that there is an asymmetry in the previous paragraph, because we privilege the unstable direction over the stable one. This is because our notion of physical measure has to do with asymptotic time averages as $n \rightarrow +\infty$. If we would instead consider the asymptotics as $n \rightarrow -\infty$, then the roles of stable and unstable objects would be reversed. We should also stress an important difference between the case when the invariant set Λ is an attractor (as in the second example) and the case when it is a Cantor set (as in the third example). In the former case the trajectories that start near Λ exhibit chaotic behavior for all time $t > 0$ (the phenomenon known as *persistent chaos*), while in the latter case the chaotic behavior occurs for a limited period of time whenever the trajectory passes by in a vicinity of Λ (the phenomenon known as *intermittent chaos*).

3. EQUILIBRIUM STATES AND THEIR RELATIVES

3.1. Hyperbolic sets. Now we make our discussion more precise and more general. We consider a smooth Riemannian manifold M and a $C^{1+\alpha}$ diffeomorphism $f: M \rightarrow M$ and restrict our attention to the dynamics of f on a *locally maximal hyperbolic set*. We recall here the basic definition and most relevant properties, referring the reader to the book of Katok and Hasselblatt [KH95, Chapter 6] for a more complete account. In what follows it is useful to keep in mind the three examples discussed above.

A *hyperbolic set* for f is a compact set $\Lambda \subset M$ with $f(\Lambda) = \Lambda$ such that for every $x \in \Lambda$, the tangent space admits a decomposition $T_x M = E^s(x) \oplus E^u(x)$ with the following properties.

- (1) The splitting is Df -invariant: $Df_x(E^\sigma(x)) = E^\sigma(fx)$ for $\sigma = s, u$.
- (2) The stable subspace $E^s(x)$ is uniformly contracting and the unstable subspace $E^u(x)$ is uniformly expanding: there are constants $C \geq 1$ and $0 < \chi < 1$ such that for every $n \geq 0$ and $v^{s,u} \in E^{s,u}(x)$, we have

$$\|Df^n v^s\| \leq C\chi^n \|v^s\| \quad \text{and} \quad \|Df^{-n} v^u\| \leq C\chi^n \|v^u\|.$$

Replacing the original Riemannian metric with an *adapted metric*,⁷ we can (and will) take $C = 1$.

In the case where $\Lambda = M$, the map f is called an *Anosov diffeomorphism*.

Of course there are some diffeomorphisms that do not have any hyperbolic sets (think of isometries), but it turns out that a very large class of diffeomorphisms do, including the examples from the previous section. These examples also have the property that Λ is *locally maximal*, meaning that there is an open set $U \supset \Lambda$ for which any invariant set $\Lambda' \subset U$ is contained in Λ ; in other words, $\Lambda = \bigcap_{n \in \mathbb{Z}} f^n(U)$. In this case every C^1 -perturbation of f also has a locally maximal hyperbolic set contained in U . In particular, the set of diffeomorphisms possessing a locally maximal hyperbolic set is open in the C^1 -topology.

A number of properties follow from the definition of a hyperbolic set. First, the subspaces $E^{s,u}(x)$ depend continuously on $x \in \Lambda$. In particular, the angle between them is uniformly away from zero. In fact, since f is $C^{1+\alpha}$, the dependence on x is Hölder continuous:

$$(3.1) \quad \rho(E^{s,u}(x), E^{s,u}(y)) \leq Kd(x, y)^\beta,$$

where ρ is the Grassmannian distance between the subspaces, d is the distance in M generated by the (adapted) Riemannian metric, and $K, \beta > 0$.

Proposition 3.1 ([KH95, Theorem 6.2.3]). *The subspaces $E^{s,u}$ can be integrated locally: for every $x \in \Lambda$, there exist local stable and unstable submanifolds $V_{\text{loc}}^{s,u}(x)$ given via the graphs of $C^{1+\alpha}$ functions $\psi_x^{s,u}: B_x^{s,u}(0, \tau) \rightarrow E^{u,s}(x)$,⁸ for which we have the following:*

- (1) $V_{\text{loc}}^{s,u}(x) = \exp_x\{v + \psi_x^{s,u}(v) : v \in B_x^{s,u}(0, \tau)\}$;
- (2) $x \in V_{\text{loc}}^{s,u}(x)$ and $T_x V_{\text{loc}}^{s,u}(x) = E^{s,u}(x)$;
- (3) $f(V_{\text{loc}}^u(x)) \supset V_{\text{loc}}^u(f(x))$ and $f(V_{\text{loc}}^s(x)) \subset V_{\text{loc}}^s(f(x))$;
- (4) there is $\lambda \in (\chi, 1)$ such that $d(f(y), f(z)) \leq \lambda d(y, z)$ for all $y, z \in V_{\text{loc}}^s(x)$ and $d(f^{-1}(y), f^{-1}(z)) \leq \lambda d(y, z)$ for all $y, z \in V_{\text{loc}}^u(x)$;⁹
- (5) there is $C > 0$ such that the Hölder seminorm satisfies $|D\psi_x^{s,u}|_\alpha \leq C$ for all $x \in \Lambda$.

The number τ is the *size* of the local manifolds, and it will be fixed at a sufficiently small value to guarantee various estimates (such as the last item in the list above). Note that the properties listed above remain true if τ is decreased. The manifolds $V_{\text{loc}}^{s,u}(x)$ depend continuously on $x \in \Lambda$.

Given a hyperbolic set Λ , there is $\varepsilon > 0$ such that for every $x, y \in \Lambda$ with $d(x, y) < \varepsilon$, the intersection $V_{\text{loc}}^s(x) \cap V_{\text{loc}}^u(y)$ consists of a single point, denoted by $[x, y]$ and called the *Smale bracket* of x and y . One can show that Λ is locally

⁷This metric may not be smooth but will be at least $C^{1+\gamma}$ for some $\gamma > 0$, which is sufficient for our purposes.

⁸Here $B_x^{s,u}(0, \tau)$ is the ball in $E^{s,u}(x) \subset T_x M$ of radius τ centered at 0.

⁹This means that the local unstable manifold for f is the local stable manifold for f^{-1} .

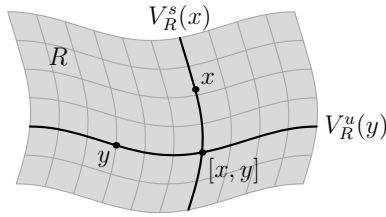


FIGURE 3.1. A rectangle in the case when $\Lambda = M$.

maximal if and only if $[x, y] \in \Lambda$ for all such x, y ; this is the local product structure referred to in §2.4.

Definition 3.2. A closed set $R \subset \Lambda$ is called a *rectangle* if $[x, y]$ is defined and lies in R for all $x, y \in R$. Given $p \in R$, we write $V_R^{s,u}(p) = V_{\text{loc}}^{s,u}(p) \cap R$ for the parts of the local manifolds that lie in R ; see Figure 3.1.

Given a rectangle R and a point $p \in R$, let $A = V_R^u(p) \subset V_{\text{loc}}^u(p)$ and $B = V_R^s(p) \subset V_{\text{loc}}^s(p)$. Then $[x, y]$ is defined for all $x \in A$ and $y \in B$, and

$$(3.2) \quad R = [A, B] := \{[x, y] : x \in A, y \in B\}.$$

Conversely, it is not hard to show that (3.2) defines a rectangle whenever $p \in \Lambda$ and the closed sets $A \subset V_{\text{loc}}^u(p) \cap \Lambda$ and $B \subset V_{\text{loc}}^s(p) \cap \Lambda$ are contained in a sufficiently small neighborhood of p .

For the hyperbolic toral automorphism from §2.1, we can take A and B to be intervals around p in the stable and unstable directions, respectively. Then $[A, B]$ is the direct product of two intervals, consistent with our usual picture of a rectangle. However, in general, we could just as easily let A and B be Cantor sets and, thus, obtain a dynamical rectangle that does not look like the picture we are familiar with. For the solenoid and horseshoe, this is the only option; in these examples the hyperbolic set Λ has zero volume and empty interior, and we see that rectangles are not even connected.

Indeed, there is a general dichotomy: given a $C^{1+\alpha}$ diffeomorphism f and a locally maximal hyperbolic set Λ , we either have $\Lambda = M$ (in which case f is an Anosov diffeomorphism) or Λ has zero volume.¹⁰ Even when $m(\Lambda) = 0$, the dynamics on Λ still influences the behavior of nearby trajectories, as is most apparent when Λ is an *attractor*, meaning that there is an open set $U \supset \Lambda$ (called a *trapping region*) such that $\overline{f(U)} \subset U$ and $\Lambda = \bigcap_{n \in \mathbb{N}} f^n(U)$, as was the case for the solenoid. In this case every trajectory that enters U is *shadowed* by some trajectory in Λ , and Λ is a union of unstable manifolds: $V_{\text{loc}}^u(x) \subset \Lambda$ for every $x \in \Lambda$.

One final comment on the topological dynamics of hyperbolic sets is in order. Recall that if X is a compact metric space and $f: X \rightarrow X$ is continuous, then the system (X, f) is called *topologically transitive* if for every open set $U, V \subset X$ there is $n \in \mathbb{N}$ such that $U \cap f^{-n}(V) \neq \emptyset$ and is called *topologically mixing* if for every such U, V there is $N \in \mathbb{N}$ such that $U \cap f^{-n}(V) \neq \emptyset$ for all $n \geq N$. Every locally maximal hyperbolic set Λ admits a *spectral decomposition* [Sma67]. It can be written as a union of disjoint closed invariant subsets $\Lambda_1, \dots, \Lambda_k \subset \Lambda$ such that each $f|_{\Lambda_i}$ is topologically transitive, and moreover each Λ_i is a union of disjoint

¹⁰This dichotomy can fail if f is only C^1 ; see [Bow75b].

closed invariant subsets $\Lambda_{i,1}, \dots, \Lambda_{i,n_i}$ such that $f(\Lambda_{i,j}) = \Lambda_{i,j+1}$ for $1 \leq j < n_i$, $f(\Lambda_{i,n_i}) = \Lambda_{i,1}$, and each $f^{n_i}|_{\Lambda_{i,j}}$ is topologically mixing. For this reason there is no real loss of generality in restricting our attention to topologically mixing locally maximal hyperbolic sets.

3.2. Conditional measures. Now we consider measures on Λ , writing $\mathcal{M}(\Lambda)$ for the set of all Borel probability measures on Λ , and $\mathcal{M}(f, \Lambda)$ for the set of all such measures that are f -invariant. We briefly mention several basic facts that play an important role in the proofs (see [EW11, Chapter 4] for details): the set $\mathcal{M}(f, \Lambda)$ is convex, and its extreme points are precisely the ergodic measures; every $\mu \in \mathcal{M}(f, \Lambda)$ has a unique *ergodic decomposition* $\mu = \int_{\mathcal{M}^e(f, \Lambda)} \nu d\zeta(\nu)$, where ζ is a probability measure on the space of ergodic measures $\mathcal{M}^e(f, \Lambda)$; and finally, $\mathcal{M}(f, \Lambda)$ is compact in the weak* topology.

As suggested by the discussion in §2.4, in order to understand how an invariant measure μ governs the forward asymptotic behavior of trajectories, we should study how μ behaves *along the unstable direction*. To make this precise, we now recall the notion of *conditional measures*; for more details, see [Roh52] or [EW11, §5.3].

Given $\mu \in \mathcal{M}(f, \Lambda)$, consider a rectangle $R \subset \Lambda$ with $\mu(R) > 0$. Let ξ be the partition of R by local unstable sets $V_R^u(x) = V_{\text{loc}}^u(x) \cap R$, $x \in R$. These depend continuously on the point x , so the partition ξ is *measurable*. This implies that the measure μ can be disintegrated with respect to ξ : for μ -almost every $x \in R$, there is a *conditional measure* $\mu_{V_R^u(x)}^\xi$ on the partition element $V_R^u(x)$ such that for any Borel subset $E \subset R$, we have¹¹

$$(3.3) \quad \mu(E) = \int_R \int_{V_R^u(x)} \mathbf{1}_E(y) d\mu_{V_R^u(x)}^\xi(y) d\mu(x).$$

Since $\mu_{V_R^u(x)}^\xi = \mu_{V_R^u(x')}^\xi$, whenever $x' \in V_R^u(x)$, the outer integral in (3.3) can also be written as an integral over the quotient space R/ξ , which inherits a factor measure $\tilde{\mu}$ from $\mu|_R$ in the natural way. By the local product structure of R , we can also fix $p \in R$ and identify R/ξ with $V_R^s(p)$. Then $\tilde{\mu}$ gives a measure on $V_R^s(p)$ by $\tilde{\mu}(A) = \mu(\bigcup_{x \in A} V_R^u(x))$. Writing $\mu_x^u = \mu_{V_R^u(x)}^\xi$ for the conditional measure on the leaf through x , we can rewrite (3.3) as¹²

$$(3.4) \quad \mu(E) = \int_{V_R^s(p)} \int_{V_R^u(x)} \mathbf{1}_E(y) d\mu_x^u(y) d\tilde{\mu}(x).$$

This disintegration is unique under the assumption that the conditional measures are normalized. Although the definition depends on R , in fact choosing a different rectangle R' merely has the effect of multiplying μ_x^u by a constant factor on $R \cap R'$ [CPZ18, Lemma 2.5].

One can similarly define a system of conditional measures $\{\mu_x^s\}$ on $V_R^s(x)$ for $x \in V_R^u(p)$. It is then natural to ask whether the conditional measure μ_p^s agrees with the measure $\tilde{\mu}$ on $V_R^s(p)$ from (3.4), and we will return to this question in §4.2.2 when we discuss *absolute continuity* and the *Hopf argument*.

¹¹For a finite partition, the obvious way to define a conditional measure on a partition element A with $\mu(A) > 0$ is to put $\mu_A(E) = \mu(E \cap A)/\mu(A)$. Roughly speaking, measurability of ξ guarantees that it can be written as a limit of finite partitions, and the conditional measures in (3.3) are the limits of the conditional measures for the finite partitions; see Proposition 6.7 for a precise statement.

¹²Note that μ_x^u depends on R , although this is suppressed in the notation.

3.3. SRB measures. Now suppose that Λ is a hyperbolic attractor and, hence, contains the local unstable leaf $V_{\text{loc}}^u(x)$ for every $x \in \Lambda$.

Definition 3.3. Given a hyperbolic attractor Λ for f and a point $y \in \Lambda$ with a local unstable leaf $W = V_{\text{loc}}^u(y)$, let m_W be the leaf volume on W generated by the restriction of the Riemannian metric to W . An invariant measure μ is a *Sinai–Ruelle–Bowen (SRB) measure* if for every rectangle $R \subset \Lambda$ with $\mu(R) > 0$, the conditional measures μ_x^u are absolutely continuous with respect to the leaf volumes m_x^u for $\tilde{\mu}$ -almost every x .

One of the major goals in the study of systems with some hyperbolicity is to construct SRB measures. In the uniformly hyperbolic setting, this was done by Sinai, Ruelle, and Bowen.

Theorem 3.4 ([Sin68, Bow75a, Rue76]). *Let Λ be a topologically transitive hyperbolic attractor for a $C^{1+\alpha}$ diffeomorphism f . Then there is a unique SRB measure for $f|_\Lambda$.*

As suggested by the discussion in §2, it is not hard to show that SRB measures are *physical* in the sense of Definition 2.1. In fact, one can prove that for hyperbolic attractors, SRB measures are the *only* physical measures.¹³

In addition to this physicality property, it was shown in [Sin68, Rue76] that the SRB measure μ has the property that

$$(3.5) \quad \mu = \lim_{n \rightarrow \infty} \frac{1}{n} \sum_{k=0}^{n-1} f_*^k m|_U,$$

where $m|_U$ is normalized volume on the trapping region $U \supset \Lambda$.¹⁴ In [PS82], this idea was used in order to *construct* SRB measures¹⁵ with $m|_U$ replaced by leaf volume $m_{V_{\text{loc}}^u(x)}$. We refer to this as the *geometric construction* of SRB measures and will return to it when we discuss our main results. First, though, we observe that the original constructions of SRB measures followed a different approach and used mathematical tools borrowed from statistical physics, as we discuss in the next section.

3.4. Equilibrium states. It turns out that it is possible to relate the absolute continuity requirement in Definition 3.3 to a variational problem. The Margulis–Ruelle inequality [Rue78a] (see also [BP13, §9.3.2]) states that for any invariant Borel measure μ supported on a hyperbolic set Λ ,¹⁶ we have the following upper bound for the Kolmogorov–Sinai entropy:

$$(3.6) \quad h_\mu(f) \leq \int \log |\det Df|_{E^u(x)}| d\mu.$$

¹³An example due to Bowen and Katok [Kat80, §0.3] shows that when Λ is not an attractor, it can support a physical measure that is not an SRB measure.

¹⁴In fact, they proved the stronger property that $f_*^n m|_U \rightarrow \mu$.

¹⁵More precisely, [PS82] considered the partially hyperbolic setting and used this approach to construct invariant measures that are absolutely continuous along unstable leaves; SRB measures are a special case of this when the center bundle is trivial.

¹⁶There is a more general version of this inequality that holds without the assumption that μ is supported on a hyperbolic set, but it requires the notion of *Lyapunov exponents*, which are beyond the scope of this paper.

Recall that $h_\mu(f)$ can be interpreted as the average asymptotic rate at which information is gained if we observe a stochastic process distributed according to μ ; (3.6) says that this rate can never exceed the average rate of expansion in the unstable direction.

Pesin's entropy formula [Pes77] states that equality holds in (3.6) if μ is absolutely continuous with respect to volume. In fact, Ledrappier and Strelcyn proved that it is sufficient for μ to have conditional measures on local unstable manifolds that are absolutely continuous with respect to leaf volume [LS82], and Ledrappier proved that this condition is also necessary [Led84]. In other words, equality holds in (3.6) if and only if μ is an SRB measure.

Since every hyperbolic attractor Λ has an SRB measure, we conclude that the function $\varphi^{\text{geo}}(x) = -\log |\det Df|_{E^u(x)}|$ has the property that

$$\sup_{\mu \in \mathcal{M}(f, \Lambda)} \left(h_\mu(f) + \int \varphi^{\text{geo}} d\mu \right) = 0,$$

and the SRB measure for f is the unique measure achieving the supremum, as claimed in §1.2. More generally, we have the following definition.

Definition 3.5. Let $\varphi: M \rightarrow \mathbb{R}$ be a continuous function, which we call a *potential*. An *equilibrium state* (or *equilibrium measure*) for φ is a measure μ achieving the supremum

$$(3.7) \quad \sup_{\mu \in \mathcal{M}(f, \Lambda)} \left(h_\mu(f) + \int \varphi d\mu \right).$$

Thus SRB measures are equilibrium states for the *geometric potential* φ^{geo} , which is Hölder continuous on every hyperbolic set as long as f is $C^{1+\alpha}$, by (3.1). This means that existence and uniqueness of SRB measures is a special case of the following classical result.

Theorem 3.6 ([Sin72, Bow75a, Rue78b]). *Let Λ be a locally maximal hyperbolic set for a $C^{1+\alpha}$ diffeomorphism f , and let $\varphi: \Lambda \rightarrow \mathbb{R}$ be a Hölder continuous potential. Assume that $f|_\Lambda$ is topologically transitive. Then there exists a unique equilibrium state μ for φ .*

In §§3.5–3.6 we briefly recall two classical proofs of Theorem 3.6 which are based on either symbolic representation of $f|_\Lambda$ as a topological Markov chain or on the specification property of $f|_\Lambda$. In §4.3 we introduce the tools that we will use to provide a new proof which is based on some constructions in geometric measure theory.

The function $\mu \mapsto h_\mu(f) + \int \varphi d\mu$ is affine. It follows that the unique equilibrium state μ must be ergodic; otherwise, every element of its ergodic decomposition would also be an equilibrium state. In fact, it has many *good* ergodic properties: one can prove that it is Bernoulli, has exponential decay of correlations, and satisfies the central limit theorem [Bow75a].

The fundamental result of thermodynamic formalism is the *variational principle*, which establishes that the supremum in (3.7) is equal to the *topological pressure* of φ , which can be defined as follows without reference to invariant measures.

Definition 3.7. Given an integer $n \geq 0$, consider the *dynamical metric of order n*

$$(3.8) \quad d_n(x, y) = \max\{d(f^k x, f^k y) : 0 \leq k < n\}$$

and the associated *Bowen balls* $B_n(x, r) = \{y : d_n(x, y) < r\}$ for each $r > 0$. We say that $E \subset \Lambda$ is (n, r) -separated if $d_n(x, y) \geq r$ for all $x \neq y \in E$, and that E is (n, r) -spanning for $X \subset \Lambda$ if $X \subset \bigcup_{x \in E} B_n(x, r)$.

Writing $S_n\varphi(x) = \sum_{k=0}^{n-1} \varphi(f^k x)$ for the n th *Birkhoff sum* along the orbit of x , the *partition sum* of φ on a set $X \subset \Lambda$ refers to one of the following two quantities:

$$Z_n^{\text{span}}(X, \varphi, r) := \inf \left\{ \sum_{x \in E} e^{S_n\varphi(x)} : E \subset X \text{ is } (n, r)\text{-spanning for } X \right\},$$

$$Z_n^{\text{sep}}(X, \varphi, r) := \sup \left\{ \sum_{x \in E} e^{S_n\varphi(x)} : E \subset X \text{ is } (n, r)\text{-separated} \right\}.$$

Then the topological pressure is given by¹⁷

$$(3.9) \quad P(\varphi) = \lim_{r \rightarrow 0} \overline{\lim}_{n \rightarrow \infty} \frac{1}{n} \log Z_n^{\text{span}}(\Lambda, \varphi, r) = \lim_{r \rightarrow 0} \overline{\lim}_{n \rightarrow \infty} \frac{1}{n} \log Z_n^{\text{sep}}(\Lambda, \varphi, r).$$

(One gets the same value if $\overline{\lim}$ is replaced by $\underline{\lim}$.)

It is worth noting at this point that the definition of $P(\varphi)$ bears a certain similarity to the definition of box dimension: one covers X by a collection of balls at a given scale, associates a certain weight to this collection, and then computes the growth rate of this weight as the balls in the cover are refined. The difference is that here the refinement is done dynamically rather than statically, and different balls carry different weight according to the ergodic sum $S_n\varphi(x)$; we will discuss this point further in §4.3 and §5. When $\varphi = 0$, we obtain the *topological entropy* $h_{\text{top}}(f) = P(0)$, which gives the asymptotic growth rate of the cardinality of an (n, r) -spanning or (n, r) -separated set. One can show that this is also the asymptotic growth rate of the number of periodic orbits in Λ of length n .

Now the variational principle [Wal82, Theorem 9.10] can be stated as follows:

$$(3.10) \quad P(\varphi) = \sup_{\mu \in \mathcal{M}(f, \Lambda)} \left(h_\mu(f) + \int \varphi d\mu \right).$$

The discussion at the beginning of this section shows that $P(\varphi^{\text{geo}}) = 0$. Given a potential φ , we see that an equilibrium state for φ is an invariant measure μ_φ such that $P(\varphi) = h_{\mu_\varphi}(f) + \int \varphi d\mu_\varphi$. For the potential function $\varphi = 0$, the equilibrium state $\mu_\varphi = \mu_0$ is the *measure of maximal entropy*.

3.5. First proof of Theorem 3.6: symbolic representation of $f|_\Lambda$. The original proof of Theorem 3.6 uses a symbolic coding of the dynamics on Λ . If $\Lambda = X_1 \cup \dots \cup X_p$, then we say that a bi-infinite sequence $\omega \in \Omega_p := \{1, \dots, p\}^{\mathbb{Z}}$ codes the orbit of $x \in \Lambda$ if $f^n x \in X_{\omega_n}$ for all $n \in \mathbb{Z}$. When Λ is a locally maximal hyperbolic set, it is not hard to show that every $\omega \in \Omega_p$ codes the orbit of at most one $x \in \Lambda$; if such an x exists, call it $\pi(\omega)$. Let $\Sigma \subset \Omega_p$ be the set of all sequences that code the orbit of some $x \in \Lambda$. Then Σ is invariant under the shift map $\sigma : \Omega_p \rightarrow \Omega_p$ defined by $(\sigma\omega)_n = \omega_{n+1}$, and the map $\pi : \Sigma \rightarrow \Lambda$ is a *topological*

¹⁷The fact that the limits coincide is given by an elementary argument comparing Z_n^{span} and Z_n^{sep} . In fact, the limit in r can be removed due to expansivity of $f|_\Lambda$; see Definition 3.9 and [Wal82, Theorem 9.6].

semiconjugacy, meaning that the following diagram commutes.

$$(3.11) \quad \begin{array}{ccc} \Sigma & \xrightarrow{\sigma} & \Sigma \\ \downarrow \pi & & \downarrow \pi \\ \Lambda & \xrightarrow{f} & \Lambda \end{array}$$

If the sets X_i overlap, then the coding map may fail to be injective. One would like to produce a coding space Σ with a nice structure for which the failure of injectivity is small. This was accomplished by Sinaĭ when $\Lambda = M$ [Sin68] and by Bowen in the general setting [Bow70, Bow75a]. They showed that things can be arranged so that Σ is defined by a nearest-neighbor condition, with the failure of injectivity confined to sets that are invisible from the point of view of equilibrium states.

Theorem 3.8 ([Bow75a]). *If Λ is a locally maximal hyperbolic set for a diffeomorphism f , then there is a Markov partition $\Lambda = X_1 \cup \dots \cup X_p$ such that each X_i is a rectangle that is the closure of its interior (in the induced topology on Λ) and the corresponding coding space Σ is a topological Markov chain*

$$(3.12) \quad \Sigma = \{\omega \in \Omega_p : f(\text{int } X_{\omega_n}) \cap \text{int } X_{\omega_{n+1}} \neq \emptyset \text{ for all } n \in \mathbb{Z}\},$$

and there is a set $\Lambda' \subset \Lambda$ such that

- (1) every $x \in \Lambda'$ has a unique preimage under π , and
- (2) if $\mu \in \mathcal{M}(f, \Lambda)$ is an equilibrium state for a Hölder continuous potential $\varphi: \Lambda \rightarrow \mathbb{R}$, then $\mu(\Lambda') = 1$.

With this result in hand, the problem of existence and uniqueness of equilibrium states can be transferred from the smooth system (Λ, f, φ) to the symbolic system $(\Sigma, \sigma, \tilde{\varphi} := \varphi \circ \pi)$, where tools from statistical mechanics and Gibbs distributions can be used. We recall here the most important ideas, referring to [Sin72, Bow75a, Rue78b] for full details.

Give Σ the metric $d(\omega, \omega') = 2^{-\min\{n: \omega_n \neq \omega'_n\}}$, so that two sequences are close if they agree on a long interval of integers around the origin. The coding map π is Hölder continuous in this metric, so $\tilde{\varphi}$ is also Hölder continuous. Fixing $r \in (\frac{1}{2}, 1)$, the Bowen balls associated to the dynamical metric (3.8) are given by

$$(3.13) \quad B_n(\omega, r) = \{\omega' \in \Sigma : \omega'_i = \omega_i \text{ for all } 0 \leq i < n\} =: C_n(\omega),$$

which we call the n -cylinder of ω . Let $E_n \subset \Sigma$ contain exactly one point from each n -cylinder. Then E_n is both (n, r) -spanning and (n, r) -separated, and writing $Z_n(\Sigma, \tilde{\varphi}) = \sum_{\omega \in E_n} e^{S_n \tilde{\varphi}(\omega)}$, one obtains

$$P(\varphi) = P(\tilde{\varphi}) = \lim_{n \rightarrow \infty} \frac{1}{n} \log Z_n(\Sigma, \tilde{\varphi}).$$

To understand what an equilibrium state for $\tilde{\varphi}$ should look like, recall that the Kolmogorov–Sinaĭ entropy of a σ -invariant measure $\tilde{\mu}$ is defined as

$$h_{\tilde{\mu}}(\sigma) = \lim_{n \rightarrow \infty} \frac{1}{n} \sum_{\omega \in E_n} -\tilde{\mu}(C_n(\omega)) \log \tilde{\mu}(C_n(\omega)).$$

A short exercise using the invariance of $\tilde{\mu}$ and the continuity of φ shows that

$$\int \tilde{\varphi} d\tilde{\mu} = \lim_{n \rightarrow \infty} \sum_{\omega \in E_n} \int_{C_n(\omega)} \tilde{\varphi} d\tilde{\mu} = \lim_{n \rightarrow \infty} \sum_{\omega \in E_n} \tilde{\mu}(C_n(\omega)) \cdot \frac{1}{n} S_n \tilde{\varphi}(\omega).$$

Thus maximizing $h_{\tilde{\mu}}(\sigma) + \int \tilde{\varphi} d\tilde{\mu}$ involves maximizing the limit of a sequence of expressions of the form $F(p_1, \dots, p_N) = \sum_{i=1}^N p_i(-\log p_i + a_i)$, where $N(n) = \#E_n$ and p_i, a_i are given by $\tilde{\mu}(C_n(\omega))$ and $S_n\tilde{\varphi}(\omega)$, so that $p_i \geq 0$ and $\sum_i p_i = 1$. It is a calculus exercise to show that with a_i fixed, F achieves its maximum value of $F = \log \sum e^{a_j} = \log Z_n(\Sigma, \tilde{\varphi}) \approx nP(\varphi)$ when $p_i = e^{a_i} / \sum_j e^{a_j} \approx e^{a_i} e^{-nP(\varphi)}$.

This last relation can be rewritten as $\tilde{\mu}(C_n(\omega)) \approx e^{S_n\tilde{\varphi}(\omega)} e^{-nP(\varphi)}$. With this in mind, one can use tools from functional analysis and statistical mechanics to show that there is a σ -invariant ergodic measure $\tilde{\mu}$ on Σ which has the *Gibbs property* with respect to $\tilde{\varphi}$: there is $Q > 0$ such that for every $\omega \in \Sigma$ and $n \in \mathbb{N}$, we have

$$(3.14) \quad Q^{-1} \leq \frac{\tilde{\mu}(C_n(\omega))}{\exp(-P(\varphi)n + S_n\tilde{\varphi}(\omega))} \leq Q.$$

By a general result that we will state momentarily, this is enough to guarantee that $\tilde{\mu}$ is the unique equilibrium state for $(\Sigma, \sigma, \tilde{\varphi})$, and hence by Theorem 3.8, its projection $\mu = \pi_*\tilde{\mu}$ is the unique equilibrium state for (Λ, f, φ) .

To formulate the link between the Gibbs property and equilibrium states, we first recall the following more general definitions.

Definition 3.9. Given a compact metric space X , a homeomorphism $f: X \rightarrow X$ is said to be *expansive* if there is $\varepsilon > 0$ such that every $x \neq y \in X$ has $d(f^n x, f^n y) > \varepsilon$ for some $n \in \mathbb{Z}$.

Definition 3.10. A measure μ on X is a *Gibbs measure* for $\varphi: X \rightarrow \mathbb{R}$ if for every small $r > 0$ there is $Q = Q(r) > 0$ such that for every $x \in X$ and $n \in \mathbb{N}$, we have

$$(3.15) \quad Q^{-1} \leq \frac{\mu(B_n(x, r))}{\exp(-P(\varphi)n + S_n\varphi(x))} \leq Q.$$

Note that $\sigma: \Sigma \rightarrow \Sigma$ is expansive, and that (3.14) implies (3.15) in this symbolic setting. Then uniqueness of the equilibrium state is a consequence of the following general result.

Proposition 3.11 ([Bow75, Lemma 8]). *If X is a compact metric space, $f: X \rightarrow X$ is an expansive homeomorphism, and μ is an ergodic f -invariant Gibbs measure for $\varphi: X \rightarrow \mathbb{R}$, then μ is the unique equilibrium state for φ .*

We remark that (3.15) does not require the Gibbs measure to be invariant. Indeed, one can separate the problem of finding a unique equilibrium state into two parts: first construct a Gibbs measure without worrying about whether or not it is invariant; then find a density function (bounded away from 0 and ∞) that produces an ergodic invariant Gibbs measure, which is the unique equilibrium state by Proposition 3.11.

3.6. Second proof of Theorem 3.6: specification property. There is another proof of Theorem 3.6 which is due to Bowen [Bow75] and avoids symbolic dynamics. Instead, it uses the fact that f satisfies the following *specification property* on a topologically mixing locally maximal hyperbolic set Λ . For each $\delta > 0$, there is an integer $p(\delta)$ such that given any points $x_1, \dots, x_n \in \Lambda$ and intervals of integers $I_1, \dots, I_n \subset [a, b]$ with $d(I_i, I_j) \geq p(\delta)$ for $i \neq j$, there is a point $x \in \Lambda$ with $f^{b-a+p(\delta)}(x) = x$ and $d(f^k(x), f^k(x_i)) < \delta$ for $k \in I_i$. Roughly speaking, f satisfies specification if for every finite number of orbit segments one can find a single periodic orbit that consecutively approximates each segment with a fixed precision

$\delta > 0$ and such that transition times are bounded by $p(\delta)$. This property allows one to study some topological and statistical properties of f by only analyzing periodic orbits.

The construction of the Gibbs measure $\tilde{\mu}$ in the first approach uses eigendata of a certain linear operator acting on an appropriately chosen Banach space of functions on Λ . The specification property allows one to use a more elementary construction and obtain a Gibbs measure on Λ as a weak* limit point of measures supported on periodic orbits. Let $\text{Per}_n := \{x \in \Lambda : f^n x = x\}$ and $Z_n^{\text{per}}(\varphi) := \sum_{x \in \text{Per}_n} e^{S_n \varphi(x)}$ (compare this to Z_n^{span} and Z_n^{sep} from Definition 3.7). Then consider the f -invariant Borel probability measures given by

$$\mu_n := \frac{1}{Z_n^{\text{per}}(\varphi)} \sum_{x \in \text{Per}_n} e^{S_n \varphi(x)} \delta_x,$$

where δ_x is the atomic probability measure with $\delta_x(\{x\}) = 1$.

Using some counting estimates on the partition sums $Z_n^{\text{per}}(\varphi)$ provided by the specification property, one can prove that every weak* limit point μ of the sequence μ_n is an ergodic Gibbs measure as in (3.15). Then Proposition 3.11 shows that μ is the unique equilibrium state for φ ; a posteriori, the sequence μ_n converges.

4. DESCRIPTION OF REFERENCE MEASURES AND MAIN RESULTS

In this section, and especially in Theorem 4.11, we describe a new proof of Theorem 3.6 that avoids Markov partitions and the specification property and instead mimics the geometric construction of SRB measures in §3.3. Given a locally maximal hyperbolic set Λ and a Hölder continuous potential $\varphi: \Lambda \rightarrow \mathbb{R}$, we define for each $x \in \Lambda$ a measure m_x^C on $X = V_{\text{loc}}^u(x) \cap \Lambda$ such that the sequence of measures

$$(4.1) \quad \mu_n = \frac{1}{n} \sum_{k=0}^{n-1} \frac{f^k m_x^C}{m_x^C(X)}$$

converges to the unique equilibrium state.¹⁸ In §4.2 we give some motivation for the properties we require the reference measures m_x^C to have, then in §4.3 we explain our construction of these measures. In §4.4 we state our main results establishing the properties of m_x^C , including how these measures can be used to prove Theorem 3.6. In §6 we outline the proofs of these results, referring to [CPZ18] for complete details and for proofs of various technical lemmas.

4.1. Conditional measures as reference measures. We start with the observation that if we were already in possession of the equilibrium state μ , then the conditional measures of μ would immediately define reference measures for which the construction just described produces μ . Indeed, suppose X is a compact topological space, $f: X \rightarrow X$ is a continuous map, and μ is a finite f -invariant ergodic Borel probability measure on X . Given $Y \subset X$ with $\mu(Y) > 0$ and a measurable partition ξ of Y , let $\tilde{\mu}$ be the corresponding factor measure on $\tilde{Y} := Y/\xi$, and let $\{\mu_W^\xi : W \in \xi\}$ be the conditional measures on partition elements.¹⁹ We prove the

¹⁸To be more precise we need first to extend m_x^C from X to a measure on Λ by assigning to any Borel set $E \subset \Lambda$ the value $m_x^C(E \cap X)$. We shall always assume that in (4.1) m_x^C is extended in this way.

¹⁹Note that ξ is not assumed to have any dynamical significance; in particular it need not be a partition into local unstable leaves, although this is the most relevant partition for our purposes.

following result in §6.4; it follows from an even more general result in ergodic theory that we state below as Proposition 6.11.

Theorem 4.1. *For $\tilde{\mu}$ -almost every $W \in \xi$, any probability measure ν on W such that $\nu \ll \mu_W^\xi$ has the property that $\nu_n := \frac{1}{n} \sum_{k=0}^{n-1} f_*^k \nu$ converges in the weak* topology to the measure μ .*

Of course, Theorem 4.1 is not much help in finding the equilibrium state μ , because we need to know μ to obtain the conditional measures μ_W^ξ . We must construct the reference measure m_x^C independently, without using any knowledge of existence of equilibrium states. Once we have done this, we will eventually show that m_x^C is equivalent to the conditional measure of the constructed equilibrium state, so our approach not only allows us to develop a new way of constructing equilibrium states, but also describes their conditional measures.

4.2. Conditions to be satisfied by reference measures. To motivate the properties that our reference measures must have, we first consider the specific case when Λ is an attractor and outline the steps in constructing SRB measures.

- (1) Given a local unstable leaf $W = V_{\text{loc}}^u(x)$ through $x \in \Lambda$ and $n \geq 0$, the image $W_n := f^n(W)$ is contained in the union of local leaves $V_{\text{loc}}^u(y_i)$ for some points $y_1, \dots, y_s \in W_n$, and the leaf volume m_x^u is pushed forward to a measure $f_*^n m_x^u$ such that $(f_*^n m_x^u)|_{V_{\text{loc}}^u(y_i)} \ll m_{y_i}^u$ for each i .
- (2) Each $\mu_n := \frac{1}{n} \sum_{k=0}^{n-1} f_*^k m_x^u$ can be written as a convex combination of measures with the form $\rho_n^i dm_{y_i}^u$ for some functions $\rho_n^i : V_{\text{loc}}^u(y_i) \rightarrow [0, \infty)$ that are uniformly bounded away from 0 and ∞ .
- (3) To show that any limit measure $\mu = \lim_{j \rightarrow \infty} \mu_{n_j}$ has absolutely continuous conditional measures on unstable leaves, first observe that given a rectangle R , the partition ξ into local unstable leaves can be approximated by a refining sequence of finite partitions ξ_ℓ , and the conditional measures μ_x^ξ are the weak* limits of the conditional measures $\mu_x^{\xi_\ell}$ as $\ell \rightarrow \infty$.
- (4) The bounds on the density functions ρ_n^i allow us to control the conditional measures $\mu_x^{\xi_\ell}$, and hence to control μ_x^ξ as well; in particular, these measures are absolutely continuous with respect to leaf volume, and thus μ is an SRB measure.

Now we describe two crucial properties of the leaf volumes m_x^u , which we will eventually need to mimic with our reference measures m_x^C . The first of these already appeared in (1.1) and describes how m_x^u scales under iteration by f ; this will let us conclude that the SRB measure μ is an equilibrium state for φ^{geo} . The second property describes how m_x^u behaves when we *slide along stable leaves* via a *holonomy map*; this issue has so far been ignored in our discussion, but plays a key role in the proof that the SRB measure μ is ergodic, and hence is the *unique* equilibrium state for φ^{geo} .

4.2.1. Scaling under iteration. Given any $x, y \in \Lambda$ and $A \subset f(V_{\text{loc}}^u(x)) \cap V_{\text{loc}}^u(y)$, we have

$$f_* m_x^u(A) = m_x^u(f^{-1}A) = \int_A |\det f^{-1}|_{E^u(z)} dm_y^u(z) = \int_A e^{\varphi^{\text{geo}}(f^{-1}z)} dm_y^u(z),$$

and so the Radon–Nikodym derivative comparing the family of measures m_x^u to their pushforwards is given in terms of the geometric potential

$$(4.2) \quad \frac{d(f_* m_x^u)}{dm_y^u}(z) = e^{\varphi^{\text{geo}}(f^{-1}(z))}.$$

Iterating this, we see that given $A \subset f^n(V_{\text{loc}}^u(x)) \cap V_{\text{loc}}^u(y)$, we have

$$(4.3) \quad f_*^n m_x^u(A) = m_x^u(f^{-n}A) = \int_A e^{\sum_{k=1}^n \varphi^{\text{geo}}(f^{-k}z)} dm_y^u(z).$$

By Hölder continuity of φ^{geo} and the fact that f^{-1} contracts uniformly along each V_{loc}^u , one can easily show that

$$(4.4) \quad \left| \sum_{k=1}^n (\varphi^{\text{geo}}(f^{-k}z_1) - \varphi^{\text{geo}}(f^{-k}z_2)) \right| \leq Q_u \text{ for all } z_1, z_2 \in V_{\text{loc}}^u(y),$$

where Q_u is a constant independent of $y \in \Lambda$, $z_1, z_2 \in V_{\text{loc}}^u(y)$, and $n \in \mathbb{N}$ (see Lemma 6.6 for details). Together with (4.3), this gives

$$(4.5) \quad e^{-Q_u} \leq \frac{m_x^u(f^{-n}A)}{e^{S_n \varphi^{\text{geo}}(x)} m_y^u(A)} \leq e^{Q_u} \text{ for all } A \subset f^n(V_{\text{loc}}^u(x)) \cap V_{\text{loc}}^u(y).$$

In particular, writing $B^u(y, r) = B(y, r) \cap V_{\text{loc}}^u(y)$, we observe that for each $r > 0$ there is a constant $K = K(r) > 0$ such that $m_y^u(B^u(y, r)) \in [K^{-1}, K]$ for all $y \in \Lambda$, and deduce from (4.5) that the u -Bowen ball

$$B_n^u(x, r) := \{z \in V_{\text{loc}}^u(x) : d_n(x, z) < r\} = f^{-n}B^u(f^n x, r)$$

admits the leaf volume estimate

$$(4.6) \quad K^{-1}e^{-Q_u} \leq \frac{m_x^u(B_n^u(x, r))}{e^{S_n \varphi^{\text{geo}}(x)}} \leq Ke^{Q_u}.$$

Definition 4.2. Consider a family of measures $\{\mu_x : x \in \Lambda\}$ such that μ_x is supported on $V_{\text{loc}}^u(x)$. We say that this family has the u -Gibbs property²⁰ with respect to the potential function $\varphi : \Lambda \rightarrow \mathbb{R}$ if there is $Q_1 = Q_1(r) > 0$ such that for all $x \in \Lambda$ and $n \in \mathbb{N}$, we have

$$(4.7) \quad Q_1^{-1} \leq \frac{\mu_x(B_n^u(x, r))}{e^{-nP(\varphi) + S_n \varphi(x)}} \leq Q_1.$$

In particular, (4.6) says that m_x^u has the u -Gibbs property with respect to the potential function φ^{geo} . Since the SRB measure μ constructed above has conditional measures that are given by multiplying the leaf volumes m_x^u by nice density functions, one can use (4.6) to ensure that the conditional measures of μ also have the u -Gibbs property. Integrating these conditional measures gives the Gibbs property for μ , and then some straightforward estimates involving $Z_n^{\text{span}}(\Lambda, \varphi, r)$ demonstrate that μ is an equilibrium state corresponding to the function φ^{geo} .

²⁰Note that this is a different notion than the idea of u -Gibbs state from [PS82].

4.2.2. *Sliding along stable leaves.* It remains, then, to show that μ is the *unique* equilibrium state for φ^{geo} ; this will follow from Proposition 3.11 if μ is proved to be ergodic. To establish ergodicity, we use the *Hopf argument*, which goes back to E. Hopf’s work on geodesic flow over surfaces [Hop39]. The first step is to observe that if μ is any invariant measure, then by Birkhoff’s ergodic theorem,²¹ for every $\psi \in L^1(\mu)$, the forward and backward ergodic averages exist and agree for μ -a.e. x :

$$(4.8) \quad \lim_{n \rightarrow \infty} \frac{1}{n} \sum_{k=0}^{n-1} \psi(f^k x) = \lim_{n \rightarrow \infty} \frac{1}{n} \sum_{k=0}^{n-1} \psi(f^{-k} x).$$

Let $\mathcal{B} \subset \Lambda$ be the set of points where the limits in (4.8) exist and agree for every continuous $\psi: \Lambda \rightarrow \mathbb{R}$; such points are called *Birkhoff regular*. For each $x \in \mathcal{B}$, write $\overline{\psi}(x)$ for the common value of these limits; note that $\overline{\psi}$ is defined μ -a.e. It is not hard to prove that μ is ergodic if and only if the function $\overline{\psi}: \mathcal{B} \rightarrow \mathbb{R}$ is constant μ -a.e. for every continuous $\psi: \Lambda \rightarrow \mathbb{R}$. By topological transitivity and the fact that $\overline{\psi} \circ f = \overline{\psi}$ on \mathcal{B} , one obtains the following standard result, whose proof we omit.

Lemma 4.3. *An f -invariant measure μ is ergodic if and only if for every continuous $\psi: \Lambda \rightarrow \mathbb{R}$ and every rectangle $R \subset \Lambda$, the function $\overline{\psi}: \mathcal{B} \cap R \rightarrow \mathbb{R}$ is constant μ -a.e.*

Now comes the central idea of the Hopf argument: Given $\psi \in C(\Lambda)$, if $\overline{\psi}(x)$ exists, then a short argument using the left-hand side of (4.8) gives $\overline{\psi}(y) = \overline{\psi}(x)$ for all $y \in \mathcal{B} \cap V_R^s(x)$. Similarly, $\overline{\psi}$ is constant on $\mathcal{B} \cap V_R^u(x)$ using the right-hand side of (4.8).

We want to conclude the proof of ergodicity by saying something like the following:

Since \mathcal{B} has full measure in R , it has full measure in almost every stable and unstable leaf in R . Thus there is $p \in R$ such that $\mathcal{B}_p := \bigcup_{x \in \mathcal{B} \cap V_R^u(p)} \mathcal{B} \cap V_R^s(x)$ has full measure in R , and by the previous paragraph, $\overline{\psi}$ is constant on \mathcal{B}_p , so Lemma 4.3 applies.

There is a subtlety involved in making this step rigorous. To begin with, the term “full measure” is used in two different ways: “ \mathcal{B} has full measure in R ” means that its complement $\mathcal{B}^c = R \setminus \mathcal{B}$ has $\mu(R \setminus \mathcal{B}) = 0$, while “ \mathcal{B} has full measure in the stable leaf $V_R^s(x)$ ” means that $\mu_x^s(\mathcal{B}^c) = 0$, where μ_x^s is the conditional measure of μ along the stable leaf. Using the analogue of (3.3)–(3.4) for the decomposition into stable leaves, we have

$$(4.9) \quad \mu(\mathcal{B}^c) = \int_R \mu_x^s(\mathcal{B}^c) d\mu(x) = \int_{V_R^u(p)} \mu_x^s(\mathcal{B}^c) d\tilde{\mu}_p(x),$$

where $\tilde{\mu}_p$ is the measure on $V_R^u(p)$ defined by

$$\tilde{\mu}_p(A) = \mu\left(\bigcup_{x \in A} V_R^s(x)\right).$$

²¹This is a more general version of the ergodic theorem than the one we mentioned in §2.1. This version applies even when μ is not ergodic, but does not require that the limits in (4.8) are equal to $\int \psi d\mu$; instead, one obtains $\int \overline{\psi} d\mu = \int \psi d\mu$, which implies the earlier version in the case when $\overline{\psi}$ is constant μ -a.e.

Let $\mathcal{B}' = \{x \in R : \mu_x^s(\mathcal{B}^c) = 0\}$. It follows that

$$0 = \mu(\mathcal{B}^c) = \int_R \mu_x^s(\mathcal{B}^c) d\mu(x) = \int_{R \setminus \mathcal{B}_0} \mu_x^s(\mathcal{B}^c) d\mu(x),$$

and since $\mu_x^s(\mathcal{B}^c) > 0$ for all $x \in R \setminus \mathcal{B}_0$ by definition, we conclude that $\mu(R \setminus \mathcal{B}') = 0$; in other words, $\mu_x^s(\mathcal{B}^c) = 0$ for μ -a.e. $x \in R$. A similar argument produces $\mathcal{B}'' \subset \mathcal{B}'$ such that $\mu(R \setminus \mathcal{B}'') = 0$ and $\mu_x^u(\mathcal{B}^c) = \mu_x^s(\mathcal{B}^c) = 0$ for every $x \in \mathcal{B}''$.

So far, things are behaving as we expect. Now, can we conclude that $\mu(\mathcal{B}_p) = \mu(R)$ for $p \in \mathcal{B}''$, thus completing the proof of ergodicity? Using (4.9), we have

$$\begin{aligned} \mu(\mathcal{B}_p) &= \int_{V_R^u(p)} \mu_x^s \left(\bigcup_{y \in \mathcal{B} \cap V_R^s(y)} \mathcal{B} \cap V_R^s(y) \right) d\tilde{\mu}_p(x) \\ &\geq \int_{\mathcal{B}'' \cap V_R^u(p)} \mu_x^s(\mathcal{B}) d\tilde{\mu}_p(x) = \tilde{\mu}_p(\mathcal{B}''). \end{aligned}$$

We would like to say that $\tilde{\mu}_p(\mathcal{B}'') = \tilde{\mu}_p(V_R^u(p)) = \mu(R)$, and conclude that \mathcal{B}_p has full μ -measure in R . We know that $\mu_p^u(\mathcal{B}'') = \mu_p^u(V_R^u(p))$, and so the proof will be complete if the answer to the following question is “yes”.

Question. Are the measures $\tilde{\mu}_p$ and μ_p^u on $V_R^u(p)$ equivalent?

Note that the measures μ_p^u are defined in terms of the foliation V_R^u , while the measures $\tilde{\mu}_p$ are defined in terms of the foliation V_R^s . We can write the measures $\tilde{\mu}_p$ in terms of μ_p^u as follows. Given $A \subset V_R^u(p)$, we have

$$(4.10) \quad \tilde{\mu}_p(A) = \mu \left(\bigcup_{x \in A} V_R^s(x) \right) = \int_R \mu_y^u \{V_R^u(y) \cap V_R^s(x) : x \in A\} d\mu(y).$$

For each $p, y \in R$, consider the (stable) holonomy map $\pi_{py} : V_R^u(p) \rightarrow V_R^u(y)$ defined by $\pi_{py}(x) = V_R^u(y) \cap V_R^s(x)$, which maps one unstable leaf to another by sliding along stable leaves; see Figure 4.1. Then (4.10) becomes

$$\tilde{\mu}_p(A) = \int_R (\mu_y^u \circ \pi_{py})(A) d\mu(y).$$

In other words, $\tilde{\mu}_p$ is the average of the conditional measures $\pi_{py}^* \mu_y^u = \mu_y^u \circ \pi_{py}$ taken over all $y \in R$.

Definition 4.4. Let $\{\nu_x\}_{x \in R}$ be a family of measures on R with the property that each ν_x is supported on $V_R^u(x)$, and let $\nu_x = \nu_y$ whenever $y \in V_R^u(x)$. We say

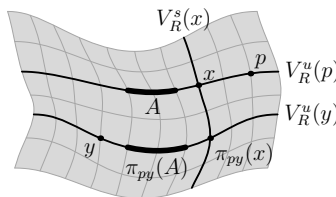


FIGURE 4.1. The stable holonomy map from $V_R^u(p)$ to $V_R^u(y)$

that the family $\{\nu_x\}$ is *absolutely continuous*²² with respect to stable holonomies if $\pi_{xy}^* \nu_y \ll \nu_x$ for all $x, y \in R$.

The preceding arguments lead to the following result; full proofs (and further discussion) can be found in [CHT16].

Proposition 4.5. *Let Λ be a topologically transitive hyperbolic set for a $C^{1+\alpha}$ diffeomorphism f , and let μ be an f -invariant measure on Λ . Suppose that for every rectangle $R \subset \Lambda$ with $\mu(R) > 0$, the unstable conditional measures μ_x^u are absolutely continuous with respect to stable holonomies. Then μ is ergodic.*

4.3. Construction of reference measures. In light of the previous section, our goal is to construct for each potential φ a reference measure m_x^C on each leaf $V_{\text{loc}}^u(x)$ satisfying a property analogous to (4.2) with φ in place of φ^{geo} , together with the absolute continuity property from Definition 4.4.

From now on we fix a local unstable manifold $W = V_{\text{loc}}^u(x)$ of size τ and consider the set $X = W \cap \Lambda$ on which we will build our reference measure. Before treating general potentials, we start with the geometric potential φ^{geo} , and we assume that Λ is an attractor for f , so that $W \subset \Lambda$. This is necessary for the moment since the measure we build will be supported on $W \cap \Lambda$, and the support of m_W is all of W ; in the general construction below we will not require Λ to be an attractor. For the geometric potential φ^{geo} , we know the reference measure m_x^C should be equivalent to the leaf volume m_W on W .²³ Leaf volume is equivalent to the Hausdorff measure $m_H(\cdot, \alpha)$ with $\alpha = \dim E^u$, which is defined by

$$(4.11) \quad m_H(Z, \alpha) := \liminf_{\varepsilon \rightarrow 0} \sum_{i=1}^{\infty} (\text{diam } U_i)^\alpha,$$

where the infimum is taken over all collections $\{U_i\}$ of open sets $U_i \subset W$ with $\text{diam } U_i \leq \varepsilon$ which cover $Z \subset W$.

We want to describe a measure that is equivalent to $m_H(\cdot, \alpha)$ but whose definition uses the dynamics of f . In (4.11), the covers used to measure Z were refined geometrically by sending $\varepsilon \rightarrow 0$. We consider instead covers that refine *dynamically*: we restrict the sets U_i to be u -Bowen balls $B_n^u(x, r) = B_n(x, r) \cap W$, and refine the covers by requiring n to be large rather than by requiring r to be small. Note that if U_i is a metric ball $B(x, \varepsilon)$, then $(\text{diam } U_i)^\alpha \approx m_W(U_i)$ up to a multiplicative factor that is bounded away from 0 and ∞ . For a u -Bowen ball, on the other hand, (4.6) gives $m_W(B_n^u(x, r)) \approx e^{S_n \varphi^{\text{geo}}(x)}$, and so we use this quantity to compute the weight of the cover. This suggests that we should fix $r > 0$ and define the measure of $Z \subset W$ by

$$(4.12) \quad m_x^{\varphi^{\text{geo}}}(Z) := \liminf_{N \rightarrow \infty} \sum_{i=1}^{\infty} e^{S_{n_i} \varphi^{\text{geo}}(x_i)},$$

where the infimum is taken over all collections $\{B_{n_i}^u(x_i, r)\}$ of u -Bowen balls with $x_i \in W$, $n_i \geq N$, which cover Z . It is relatively straightforward to derive property (4.2) from (4.12).

²²There is a related, but distinct, notion of absolute continuity of a foliation (with respect to volume), which also plays a key role in smooth ergodic theory; see [BP07, §8.6].

²³From Theorem 4.1 we see that the equivalence class of the measure is the crucial thing for the geometric construction to work.

Now it is nearly apparent what the definition should be for a general potential: we want to replace φ^{geo} with φ in (4.12). There is one small subtlety, though. First, Definition 3.7 gives $P(\varphi + c) = P(\varphi) + c$ for $c \in \mathbb{R}$. This along with the definition of equilibrium state and the variational principle (3.10) shows that adding a constant to φ does not change its equilibrium states, and thus we should also expect that φ and $\varphi + c$ produce the same reference measure on $W \cap \Lambda$. For this to happen, we need to modify (4.12) so that adding a constant to φ does not affect the value. This can be achieved by multiplying each term in the sum by $e^{-n_i P(\varphi)}$. Note that since $P(\varphi^{\text{geo}}) = 0$, this does not modify (4.12). Thus we make the following definition.

Definition 4.6. Let $0 < r < \frac{\tau}{3}$. We define a measure on $X := W \cap \Lambda$ by

$$(4.13) \quad m_x^{\mathcal{C}}(Z) := \liminf_{N \rightarrow \infty} \sum_i e^{-n_i P(\varphi)} e^{S_{n_i} \varphi(x_i)},$$

where the infimum is taken over all collections $\{B_{n_i}^u(x_i, r)\}$ of u -Bowen balls with $x_i \in X$, $n_i \geq N$, which cover Z , and for convenience we write $\mathcal{C} = (\varphi, r)$ to keep track of the data on which the reference measure depends.

Both definitions (4.12) and (4.13) are specific cases of the Carathéodory measure produced by a dynamically defined Carathéodory dimension structure, which we discuss at greater length in §5; this is the Pesin–Pitskel’ definition of topological pressure [PP84] that generalized Bowen’s definition of topological entropy for noncompact sets [Bow73]. In particular, Proposition 5.4 establishes the crucial property that every local unstable leaf carries the same topological pressure as the entire set Λ .

4.4. Statements of main results. Now we state the most important properties of $m_x^{\mathcal{C}}$ and show how it can be used as a reference measure to construct the equilibrium state for φ . All results in this section are proved in detail in [CPZ18];²⁴ we outline the proofs in §6. Our first main result shows that the measure $m_x^{\mathcal{C}}$ is finite and nonzero.

Theorem 4.7 ([CPZ18, Theorem 4.2]). *Let Λ be a topologically transitive locally maximal hyperbolic set for a $C^{1+\alpha}$ diffeomorphism f , and let $\varphi: \Lambda \rightarrow \mathbb{R}$ be Hölder continuous. Fix r as in Definition 4.6, and for each $x \in \Lambda$, let $m_x^{\mathcal{C}}$ be given by (4.13), where $\mathcal{C} = (\varphi, r)$. Then there is $K > 0$ such that for every $x \in \Lambda$, $m_x^{\mathcal{C}}$ is a Borel measure on $V_{\text{loc}}^u(x) \cap \Lambda$ with $m_x^{\mathcal{C}}(V_{\text{loc}}^u(x) \cap \Lambda) \in [K^{-1}, K]$. If $V_{\text{loc}}^u(x) \cap V_{\text{loc}}^u(y) \cap \Lambda \neq \emptyset$, then $m_x^{\mathcal{C}}$ and $m_y^{\mathcal{C}}$ agree on the intersection.*

As described in §4.2, we need to understand how the measures $m_x^{\mathcal{C}}$ transform under (1) the dynamics of f and (2) sliding along stable leaves via holonomy. For the first of these properties, the following result gives the necessary scaling property analogous to (4.2).

Theorem 4.8 ([CPZ18, Theorem 4.4]). *Under the hypotheses of Theorem 4.7, for every $x \in \Lambda$, we have $f^* m_{f(x)}^{\mathcal{C}} := m_{f(x)}^{\mathcal{C}} \circ f \ll m_x^{\mathcal{C}}$, with Radon–Nikodym derivative $e^{P(\varphi) - \varphi}$, so that (1.2) holds.*

²⁴The numbering of references within [CPZ18] refers to the first arXiv version; it is possible that the numbering will change between this and the final published version.

Corollary 4.9 ([CPZ18, Corollary 4.5]). *Under the hypotheses of Theorem 4.7, the family of measures $\{m_x^C\}_{x \in \Lambda}$ has the u -Gibbs property from Definition 4.2. In particular, for every relatively open $U \subset V_{\text{loc}}^u(x) \cap \Lambda$, we have $m_x^C(U) > 0$.*

The final crucial property of the reference measures is that they are absolutely continuous under holonomy.

Theorem 4.10 ([CPZ18, Theorem 4.6]). *Under the hypotheses of Theorem 4.7, there is a constant $Q_2 > 0$ such that for every rectangle $R \subset \Lambda$ and every $y, z \in R$, the measures $\pi_{yz}^* m_z^C = m_z^C \circ \pi_{yz}$ and m_y^C are equivalent on $V_R^u(y)$, with*

$$Q_2^{-1} \leq \frac{d\pi_{yz}^* m_z^C}{dm_y^C} \leq Q_2.$$

Note that Theorem 4.10 in particular shows that given a rectangle $R \subset \Lambda$, if $m_x^C(V_R^u(x)) > 0$ for some $x \in R$, then the same is true for every $x \in R$. Moreover, by Corollary 4.9 this happens whenever R is the closure of its interior (relative to Λ).

Using these properties of the measures m_x^C , we can carry out the geometric construction of equilibrium states; see §6 for the proof of the following.

Theorem 4.11. *Under the hypotheses of Theorem 4.7, the following are true.*

- (1) *For every $x \in \Lambda$, the sequence of measures $\mu_n := \frac{1}{n} \sum_{k=0}^{n-1} f^k m_x^C / m_x^C(V_{\text{loc}}^u(x))$ from (4.1) is weak* convergent as $n \rightarrow \infty$ to a probability measure μ_φ that is independent of x .*
- (2) *The measure μ_φ is ergodic, gives positive weight to every open set in Λ , has the Gibbs property (3.15), and is the unique equilibrium state for (Λ, f, φ) .*
- (3) *For every rectangle $R \subset \Lambda$ with $\mu_\varphi(R) > 0$, the conditional measures μ_y^u generated by μ_φ on unstable sets $V_R^u(y)$ are equivalent for μ_φ -almost every $y \in R$ to the reference measures $m_y^C|_{V_R^u(y)}$. Moreover, there exists $Q_3 > 0$, independent of R and y , such that for μ_φ -almost every $y \in R$ we have²⁵*

$$(4.14) \quad Q_3^{-1} \leq \frac{d\mu_y^u}{dm_y^C}(z) m_y^C(R) \leq Q_3 \text{ for } \mu_y^u\text{-a.e. } z \in V_R^u(y).$$

Theorems 4.10 and 4.11(3) allow us to show that the equilibrium state μ_φ has local product structure, as follows. Consider a rectangle $R \subset \Lambda$ with $\mu_\varphi(R) > 0$, and a system of conditional measures μ_x^u with respect to the partition ξ of R into local unstable leaves. Given $p \in R$, define a measure $\tilde{\mu}_p$ on $V_R^s(p)$ by $\tilde{\mu}_p(A) = \mu_\varphi(\bigcup_{x \in A} V_R^u(x))$ as in the paragraph preceding (3.4). Since R is homeomorphic to the direct product of $V_R^u(p)$ and $V_R^s(p)$, the product of the measures μ_p^u and $\tilde{\mu}_p$ gives a measure on R that we denote by $\mu_p^u \otimes \tilde{\mu}_p$. The following local product structure result is a consequence of Theorem 4.10, Theorem 4.11(3), and (3.4); see §6.3.3.

Corollary 4.12. *For every rectangle R and μ_φ -almost every $p \in R$, we have $\pi_{py}^* \mu_y^u \sim \mu_p^u$ for $\tilde{\mu}_p$ -almost every $y \in V_R^s(p)$, and thus $\mu_\varphi \sim \mu_p^u \otimes \tilde{\mu}_p$. Moreover, it follows that $\tilde{\mu}_p$ is equivalent to μ_p^s , the conditional measure on $V_R^s(p)$ with respect to the partition into stable leaves, for μ_φ -a.e. $p \in R$.*

²⁵It is reasonable to expect, based on analogy with the case of SRB measure, that the Radon–Nikodym derivative in (4.14) is in fact Hölder continuous and is given by an explicit formula. At present we can only prove this for a modified version of m_x^C , whose definition we omit here.

We remark that Corollary 4.12 was also proved by Leplaideur [Lep00]. His proof uses Markov partitions to construct families of leaf measures with the properties given in Theorems 4.7 and 4.10. Historically, this description of μ_φ in terms of its direct product structure dates back to Margulis [Mar70], who described the unique measure of maximal entropy for a transitive Anosov flow as a direct product of leaf-wise measures satisfying the continuous-time analogue of (1.2) for $\varphi = 0$. In this specific case the equivalences in Corollary 4.12 can be strengthened to equalities.

5. CARATHÉODORY DIMENSION STRUCTURE

The definition of the measures m_x^C in (4.13) is a specific instance of the *Carathéodory dimension construction* introduced by the second author in [Pes88] (see also [Pes97, §10]). It is a substantial generalization and adaptation to dynamical systems of the classical construction of Carathéodory measure in geometric measure theory, of which Lebesgue measure and Hausdorff measure are the most well-known examples. We briefly recall here the Carathéodory dimension construction together with some of its basic properties.

5.1. Carathéodory dimension and measure. A *Carathéodory dimension structure*, or *C-structure*, on a set X is given by the following data.

- (1) An indexed collection of subsets of X , denoted $\mathcal{F} = \{U_s : s \in \mathcal{S}\}$.
- (2) Functions $\xi, \eta, \psi : \mathcal{S} \rightarrow [0, \infty)$ satisfying the following conditions:
 - (A1) if $U_s = \emptyset$, then $\eta(s) = \psi(s) = 0$; if $U_s \neq \emptyset$, then $\eta(s) > 0$ and $\psi(s) > 0$;²⁶
 - (A2) for any $\delta > 0$ one can find $\varepsilon > 0$ such that $\eta(s) \leq \delta$ for any $s \in \mathcal{S}$ with $\psi(s) \leq \varepsilon$;
 - (A3) for any $\varepsilon > 0$ there exists a finite or countable subcollection $\mathcal{G} \subset \mathcal{S}$ that covers X (meaning that $\bigcup_{s \in \mathcal{G}} U_s \supset X$) and has $\psi(\mathcal{G}) := \sup\{\psi(s) : s \in \mathcal{S}\} \leq \varepsilon$.

Note that no conditions are placed on $\xi(s)$, which we interpret as the *weight* of U_s . The values $\eta(s)$ and $\psi(s)$ can each be interpreted as a *size* or *scale* of U_s ; we allow these functions to be different from each other.

The *C-structure* $(\mathcal{S}, \mathcal{F}, \xi, \eta, \psi)$ determines a one-parameter family of outer measures on X as follows. Fix a nonempty set $Z \subset X$, and consider some $\mathcal{G} \subset \mathcal{S}$ that covers Z (meaning that $\bigcup_{s \in \mathcal{G}} U_s \supset Z$). Then $\psi(\mathcal{G})$ is interpreted as the largest size of sets in the cover, and we set for each $\alpha \in \mathbb{R}$,

$$(5.1) \quad m_C(Z, \alpha) := \liminf_{\varepsilon \rightarrow 0} \inf_{\mathcal{G}} \sum_{s \in \mathcal{G}} \xi(s) \eta(s)^\alpha,$$

where the infimum is taken over all finite or countable $\mathcal{G} \subset \mathcal{S}$ covering Z with $\psi(\mathcal{G}) \leq \varepsilon$. Defining $m_C(\emptyset, \alpha) := 0$, it follows from [Pes97, Proposition 1.1] that $m_C(\cdot, \alpha)$ is an outer measure. The measure induced by $m_C(\cdot, \alpha)$ on the σ -algebra of measurable sets is the α -*Carathéodory measure*; it need not be σ -finite or nontrivial.

Proposition 5.1 ([Pes97, Proposition 1.2]). *For any set $Z \subset X$, there exists a critical value $\alpha_C \in \mathbb{R}$ such that $m_C(Z, \alpha) = \infty$ for $\alpha < \alpha_C$ and $m_C(Z, \alpha) = 0$ for $\alpha > \alpha_C$.*

²⁶In [Pes97], condition (A1) includes the requirement that there is $s_0 \in \mathcal{S}$ such that $U_{s_0} = \emptyset$. Here we remove this assumption and instead define $m_C(\emptyset, \alpha) := 0$, which is equivalent.

We call $\dim_C Z = \alpha_C$ the *Carathéodory dimension* of the set Z associated to the C -structure $(\mathcal{S}, \mathcal{F}, \xi, \eta, \psi)$. By Proposition 5.1, $\alpha = \dim_C X$ is the only value of α for which (5.1) can possibly produce a nonzero finite measure on X , though it is still possible that $m_C(X, \dim_C X)$ is equal to 0 or ∞ .

5.2. Examples of C -structures. The C -structures in which we are interested are generated by other structures on the set X .

5.2.1. Hausdorff dimension and measure. If X is a metric space, then consider the C -structure given by $\mathcal{S} := X \times (0, \infty)$ and

$$\mathcal{F} := \{B(x, r) : x \in X, r > 0\}, \quad \xi(x, r) = 1, \quad \eta(x, r) = \psi(x, r) = r.$$

Comparing (4.11) and (5.1), we see that $m_C(Z, \alpha) = m_H(Z, \alpha)$ for every $Z \subset X$, and the Hausdorff dimension $\dim_H(Z)$ is the critical value such that $m_H(Z, \alpha)$ is infinite for $\alpha < \dim_H(Z)$ and 0 for $\alpha > \dim_H(Z)$. Thus $\dim_C Z = \dim_H Z$, and the outer measure $m_C(\cdot, \dim_H Z)$ on Z is the $(\dim_H Z)$ -dimensional spherical Hausdorff measure.

It is useful to understand when an outer measure defines a Borel measure on a metric space. Recall that an outer measure m on a metric space (X, d) is a *metric outer measure* if $m(E \cup F) = m(E) + m(F)$ whenever

$$d(E, F) := \inf\{d(x, y) : x \in E, y \in F\} > 0.$$

Proposition 5.2 ([Fed69, §2.3.2(9)]). *If X is a metric space and m is a metric outer measure on X , then every Borel set in X is m -measurable, and so m defines a Borel measure on X .*

Given any $E, F \subset X$, with $d(E, F) > 0$, we see that any cover $\mathcal{G} \subset \mathcal{F}$ of $E \cup F$ with $\psi(\mathcal{G}) \leq d(E, F)/2$ can be written as the disjoint union of a cover of E and a cover of F ; using this it is easy to show that $m_H(E \cup F, \alpha) = m_H(E, \alpha) + m_H(F, \alpha)$, so $m_H(\cdot, \alpha)$ is a metric outer measure. By Proposition 5.2, this defines a Borel measure on X .

5.2.2. Topological pressure as a Carathéodory dimension. Let f be a continuous map of a compact metric space X , and let $\varphi : X \rightarrow \mathbb{R}$ be a continuous function. Then as described already in §4.3, one can consider covers that are refined dynamically rather than geometrically. This was done first by Bowen to define topological entropy in a more general setting [Bow73], and then extended by Pesin and Pitskel' to topological pressure [PP84]. Here we give a definition that differs slightly from [PP84] but gives the same dimensional quantity [Cli11, Proposition 5.2].

Fix $r > 0$ and to each $(x, n) \in X \times \mathbb{N}$, associate the Bowen ball $B_n(x, r)$. Let \mathcal{F} be the collection of all such Bowen balls, and let $\mathcal{S} = X \times \mathbb{N}$, so $s = (x, n)$ has $U_s = B_n(x, r)$. Now put

$$(5.2) \quad \xi(x, n) = e^{S_n \varphi(x)}, \quad \eta(x, n) = e^{-n}, \quad \psi(x, n) = \frac{1}{n}.$$

It is easy to see that $(\mathcal{S}, \mathcal{F}, \xi, \eta, \psi)$ satisfies (A1)–(A3), so this defines a C -structure. The associated outer measure is given by

$$(5.3) \quad m_C(Z, \alpha) = \liminf_{N \rightarrow \infty} \inf_{\mathcal{G}} \sum_{(x, n) \in \mathcal{G}} e^{S_n \varphi(x)} e^{-n\alpha},$$

where the infimum is over all $\mathcal{G} \subset \mathcal{S}$ such that $\bigcup_{(x, n) \in \mathcal{G}} B_n(x, r) \supset Z$ and $n \geq N$ for all $(x, n) \in \mathcal{G}$.

Remark 5.3. The measure $m_C(\cdot, \alpha)$ is not necessarily a metric outer measure, since there may be $x \neq y \in X$ such that $y \in B_n(x, r)$ for all $n \in \mathbb{N}$.²⁷ Thus, Borel sets in X need not be $m_C(\cdot, \alpha)$ -measurable.

Writing $\dim_C^r Z$ for the critical value of α , where the superscript emphasizes the dependence on r , the quantity

$$P_Z(\varphi) = \overline{\lim}_{r \rightarrow 0} \dim_C^r Z$$

is called the the *topological pressure* of φ on the set Z . Observe that this notion of topological pressure is more general than the one introduced in Definition 3.7 as it is more suited to arbitrary subsets Z (which need not be compact or invariant); both definitions agree when $Z = X$ [Pes97, Theorem 11.5].

5.2.3. *A C-structure on local unstable leaves.* Now consider the setting of Theorems 4.7–4.10. Λ is a hyperbolic set for a $C^{1+\alpha}$ diffeomorphism f , and $\varphi: \Lambda \rightarrow \mathbb{R}$ is Hölder continuous. Fix $r > 0$ and define a C -structure on $X = V_{\text{loc}}^u(x) \cap \Lambda$, which depends on φ , in the following way. To each $(x, n) \in X \times \mathbb{N}$, associate the Bowen ball $B_n(x, r)$. Let \mathcal{F} be the collection of all such balls, and let $\mathcal{S} = X \times \mathbb{N}$, so $s = (x, n)$ has $U_s = B_n(x, r)$. Now put

$$(5.4) \quad \xi(x, n) = e^{S_n \varphi(x)}, \quad \eta(x, n) = e^{-n}, \quad \psi(x, n) = \frac{1}{n}.$$

Again, $(\mathcal{S}, \mathcal{F}, \xi, \eta, \psi)$ satisfies (A1)–(A3) and defines a C -structure, whose associated outer measure is given by

$$(5.5) \quad m_C(Z, \alpha) = \lim_{N \rightarrow \infty} \inf_{\mathcal{G}} \sum_{(x,n) \in \mathcal{G}} e^{S_n \varphi(x)} e^{-n\alpha},$$

where the infimum is over all $\mathcal{G} \subset \mathcal{S}$ such that $\bigcup_{(x,n) \in \mathcal{G}} B_n(x, r) \supset Z$ and $n \geq N$ for all $(x, n) \in \mathcal{G}$.

Given $x \in \Lambda$, we are interested in two things:

- (1) the Carathéodory dimension of X , as determined by this C -structure; and
- (2) the (outer) measure on X defined by (5.3) at $\alpha = \dim_C(X)$.

The first of these is settled by the following, which is proved in [CPZ18, Theorem 4.2(1)].

Proposition 5.4. *With Λ, f, φ, r as above, and the C -structure defined on $X = V_{\text{loc}}^u(x) \cap \Lambda$ by Bowen balls $B_n(x, r)$ and (5.4), we have $\dim_C^r(X) = P(\varphi)$ for every $x \in \Lambda$. In particular, this implies that $P_X(\varphi) = P_\Lambda(\varphi)$.*

Note that on each $X = V_{\text{loc}}^u(x) \cap \Lambda$, covers by Bowen balls $B_n(x, r)$ are the same thing as covers by u -Bowen balls $B_n^u(x, r) = B_n(x, r) \cap V_{\text{loc}}^u(x)$, which we used in §4.3. Thus when we put $\alpha = P(\varphi)$, we see that (5.5) agrees with (4.13) for every $Z \subset X$, and in particular, the quantity $m_x^C(Z)$ defined in (4.13) is the outer measure on X associated to the C -structure above and the parameter value $\alpha = P(\varphi)$.

One must still do some work to show that this outer measure is finite and nonzero; this is done in [CPZ18], and the idea of the argument is given in §6.1 below. We conclude this section by observing that the issue raised in Remark 5.3 is not a problem here, and that we have in fact defined a *metric* outer measure. Indeed, given any $x \in \Lambda$ and $y \in V_{\text{loc}}^u(x) \cap \Lambda$, we have $\text{diam } B_n^u(y, r) \leq r\lambda^n$ for all $n \in \mathbb{N}$ by Proposition 3.1, so if $E, F \subset X$ have $d(E, F) > 0$, then there is $N \in \mathbb{N}$ such

²⁷In fact, $m_C(\cdot, \alpha)$ is an outer measure if and only if f is positively expansive to scale r .

that $B_n(y, r) \cap B_k(z, r) = \emptyset$ whenever $y \in E, z \in F$, and $k, n \geq N$. Then for N sufficiently large, any \mathcal{G} as in (5.5) has the property that it splits into disjoint covers of E and F , and thus $m_x^C(E \cup F) = m_x^C(E) + m_x^C(F)$. By Proposition 5.2, m_x^C defines a Borel measure on X , as claimed in Theorem 4.7.

5.3. An application: measures of maximal dimension. If X is a measurable space with a measure μ , and \dim_C is a Carathéodory dimension on X , then the quantity

$$\begin{aligned} \dim_C \mu &= \inf\{\dim_C Z : \mu(Z) = 1\} \\ &= \liminf_{\delta \rightarrow 0} \{\dim_C Z : \mu(Z) > 1 - \delta\} \end{aligned}$$

is called the *Carathéodory dimension* of μ . We say that μ is a *measure of maximal Carathéodory dimension* if $\dim_C \mu = \dim_C X$. Note that if the Carathéodory measure $m_C(X, \alpha)$ at dimension $\alpha = \dim_C X$ is finite and positive, then this measure is a measure of maximal Carathéodory dimension.

With $f: \Lambda \rightarrow \Lambda$ as in Theorem 4.7, we consider a particular but important family of potential functions $\varphi_t^{\text{geo}}(x)$ on Λ , called the *geometric t -potentials*: for any $t \in \mathbb{R}$

$$\varphi_t^{\text{geo}}(x) := -t \log |\det Df|_{E^u(x)}|.$$

Since the subspace $E^u(x)$ depends Hölder continuously on $x \in \Lambda$ (see (3.1)), for each $t \in \mathbb{R}$ the function $\varphi_t^{\text{geo}}(x)$ is Hölder continuous and hence, it admits a unique equilibrium state $\mu_t := \mu_{\varphi_t^{\text{geo}}}$.

We consider the function $P(t) := P(\varphi_t^{\text{geo}})$ called the *pressure function*. One can show that this function is monotonically decreasing, convex, and real analytic in t . Moreover, $P(t) \rightarrow +\infty$ as $t \rightarrow -\infty$ and $P(t) \rightarrow -\infty$ as $t \rightarrow +\infty$ with $P(1) \leq 0$. Therefore, there is a number $0 < t_0 \leq 1$ which is the unique solution of *Bowen’s equation* $P(t) = 0$. We shall show that given $x \in \Lambda$, there is a C -structure on the set $X = V_{\text{loc}}^u(x) \cap \Lambda$ with respect to which t_0 is the Carathéodory dimension of the set X . Indeed, since $P(t_0) = 0$, the measure $m_x^{t_0} := m_x^C$, given by (4.13) for $\mathcal{C} = (\varphi_{t_0}^{\text{geo}}, r)$, can be written as

$$(5.6) \quad m_x^{t_0}(Z) = \liminf_{N \rightarrow \infty} \left\{ \sum_i \left(\prod_{k=0}^{n_i-1} \det(Df|_{E^u(f^k(x_i))}) \right)^{-t_0} \right\},$$

where the infimum is taken over all collections $\{B_{n_i}^u(x_i, r)\}$ of u -Bowen balls with $x_i \in X, n_i \geq N$, which cover Z .

Relation (5.6) shows that the measure $m_x^{t_0}$ is the Carathéodory measure generated by the C -structure $\tau' = (\mathcal{S}, \mathcal{F}, \xi', \eta', \psi)$, where

$$\xi'(x, n) := 1, \quad \eta'(x, n) := \prod_{k=0}^{n-1} \det(Df|_{E^u(f^k(x_i))})^{-1}.$$

It is easy to see that with respect to the C -structure τ' we have that $\dim_{C, \tau'} X = t_0$ and the measure $m_x^{t_0} = m_{C, \tau'}(\cdot, t)$ is the measure of maximal Carathéodory dimension. In particular, the Carathéodory dimension of $X = V_{\text{loc}}^u(x) \cap \Lambda$ does not depend on the choice of the point $x \in \Lambda$. It is also clear that the number t_0 depends continuously on f in the C^1 topology and hence, so does the Carathéodory dimension $\dim_{C, \tau'} X$.

We consider the particular case when the map f is u -conformal; that is $Df|_{E^u(x)} = a(x)\text{Isom}_x$ for all $x \in \Lambda$, where Isom_x is an isometry. The direct calculation involving (5.6) shows that in this case $m_x^{t_0}$ is a measure of full Hausdorff dimension and that $t_0 \dim E^u = \dim_H X$.

Given a locally maximal hyperbolic set Λ , it has been a long-standing open problem to compute the Hausdorff dimension of the set $X = V_{\text{loc}}^u(x) \cap \Lambda$ and to find an invariant measure whose conditional measures on unstable leaves have maximal Hausdorff dimension, provided such a measure exists. The above result solves this problem for u -conformal diffeomorphisms. The reader can find the original proof and relevant references in [Pes97]. It was recently proved that without the assumption of u -conformality, there are examples for which there is no invariant measure whose conditionals have full Hausdorff dimension; see [DS17]. Theorem 4.11 provides one way to settle the issue in the nonconformal case by replacing “measure of maximal Hausdorff dimension” with “measure of maximal Carathéodory dimension” with respect to the C -structure τ' just described.

6. OUTLINE OF PROOFS

In §§6.1–6.2 we outline the proofs of Theorems 4.7–4.10, referring to [CPZ18] for complete details; see Remarks 2.3 and 4.1 of that paper for an explanation of why the setting here is covered. In §6.3 we prove Theorem 4.11, again referring to [CPZ18] for certain technicalities. In §6.4 we give a complete proof of Theorem 4.1.

6.1. Reference measures are nonzero and finite. Recall that Λ is a locally maximal hyperbolic set for f , on which each x has local stable and unstable manifolds of size $\tau > 0$. We assume that $f|_{\Lambda}$ is topologically transitive. In what follows, we occasionally use the following notation: given $A, B, C, a \geq 0$, we write $A = C^{\pm a}B$ as shorthand to mean $C^{-a}B \leq A \leq C^aB$. The key to the proof of Theorem 4.7 is the following result.

Proposition 6.1. *For every $r_1 \in (0, \tau)$ and $r_2 \in (0, \tau/3]$, there is $C > 1$ such that for every $x \in \Lambda$ and $n \in \mathbb{N}$ we have*

$$(6.1) \quad Z_n^{\text{span}}(B_{\Lambda}^u(x, r_1), \varphi, r_2) = C^{\pm 1}e^{nP(\varphi)}.$$

Similar partition sum bounds are obtained in Bowen’s paper [Bow75], where they are proved for all of Λ instead of for a single unstable leaf. For the full proof of Proposition 6.1, see [CPZ18, §6]; we outline the argument below. As in Bowen’s case, the underlying mechanism is a set of elementary lemmas, which we give in §6.1.1. In §6.1.2 we explain why it is reasonable to expect these lemmas to apply to the sequence Z_n^{span} , and in §6.1.3 we outline how Proposition 6.1 leads to Theorem 4.7.

6.1.1. Elementary counting lemmas.

Lemma 6.2. *If $Z_n > 0$ is a sequence of numbers satisfying $Z_{n+m} \leq Z_n Z_m$ for all m, n , then $P = \lim_{n \rightarrow \infty} \frac{1}{n} \log Z_n$ exists and is equal to $\inf_{n \in \mathbb{N}} \frac{1}{n} \log Z_n$. In particular, $Z_n \geq e^{nP}$ for every n .*

Proof. Fix $n \in \mathbb{N}$. Then for all $m \in \mathbb{N}$, we can write $m = an + b$, where $a \in \mathbb{N}$ and $b \in \{0, 1, \dots, n-1\}$, and iterate the submultiplicativity property to obtain

$Z_m \leq Z_n^a Z_b$. Taking logs and dividing by m gives

$$\frac{1}{m} \log Z_m \leq \frac{a}{m} \log Z_n + \frac{\log Z_b}{m} \leq \frac{an}{m} \cdot \frac{1}{n} \log Z_n + \frac{\max\{\log Z_0, \dots, \log Z_{n-1}\}}{m}.$$

Sending $m \rightarrow \infty$, we see that $\frac{an}{m} \rightarrow 1$, so

$$(6.2) \quad \overline{\lim}_{m \rightarrow \infty} \frac{1}{m} \log Z_m \leq \frac{1}{n} \log Z_n.$$

Since n was arbitrary, we deduce that

$$\overline{\lim}_{m \rightarrow \infty} \frac{1}{m} \log Z_m \leq \inf_{n \in \mathbb{N}} \frac{1}{n} \log Z_n \leq \underline{\lim}_{n \rightarrow \infty} \frac{1}{n} \log Z_n,$$

whence all three terms are equal and the limit exists. Now (6.2) implies that $Z_n \geq e^{nP}$. □

Lemma 6.3. *If $Z_n > 0$ is a sequence of numbers satisfying $Z_{n+m} \leq CZ_n Z_m$ for all m, n , where $C > 0$ is independent of m, n , then $P = \lim_{n \rightarrow \infty} \frac{1}{n} \log Z_n$ exists and is equal to $\inf_{n \in \mathbb{N}} \frac{1}{n} \log(CZ_n)$. In particular, $Z_n \geq C^{-1} e^{nP}$ for all n .*

Proof. The proof follows by applying Lemma 6.2 to the sequence $Y_n = CZ_n$, which satisfies $Y_{n+m} = CZ_{n+m} \leq C^2 Z_n Z_m = Y_n Y_m$. □

Lemma 6.4. *If $Z_n > 0$ is a sequence of numbers satisfying $Z_{n+m} \geq C^{-1} Z_n Z_m$ for all m, n , where $C > 0$ is independent of m, n , then $P = \lim_{n \rightarrow \infty} \frac{1}{n} \log Z_n$ exists and is equal to $\sup_{n \in \mathbb{N}} \frac{1}{n} \log(Z_n/C)$. In particular, $Z_n \leq C e^{nP}$ for all n .*

Proof. The proof follows by applying Lemma 6.2 to the sequence $Y_n = C/Z_n$, which satisfies $Y_{n+m} = C/Z_{n+m} \leq C^2/(Z_n Z_m) = Y_n Y_m$. □

6.1.2. *Partition sums are nearly multiplicative.* In light of Lemmas 6.3 and 6.4, Proposition 6.1 can be proved by showing that the partition sums $Z_n^{\text{span}}(B_\Lambda^u(x, r_1), \varphi, r_2)$ are *nearly multiplicative*: $Z_{n+m}^{\text{span}} = C^{\pm 1} Z_n^{\text{span}} Z_m^{\text{span}}$. A short argument given in [CPZ18, Lemma 6.3] shows that $Z_n^{\text{sep}} = e^{\pm Q_u} Z_n^{\text{span}}$, and thus it suffices to show that

$$Z_{n+m}^{\text{span}} \leq CZ_n^{\text{span}} Z_m^{\text{span}}, \quad Z_{n+m}^{\text{sep}} \geq C^{-1} Z_n^{\text{sep}} Z_m^{\text{sep}},$$

where we are being deliberately vague about the arguments of Z^{span} and Z^{sep} .

Figure 6.1 illustrates the idea driving the estimate for Z^{sep} : if $E_n = \{y_1, \dots, y_a\} \subset B_\Lambda^u(x, r_1)$ is a maximal (n, r_2) -separated set of points, and to each $1 \leq i \leq a$ we

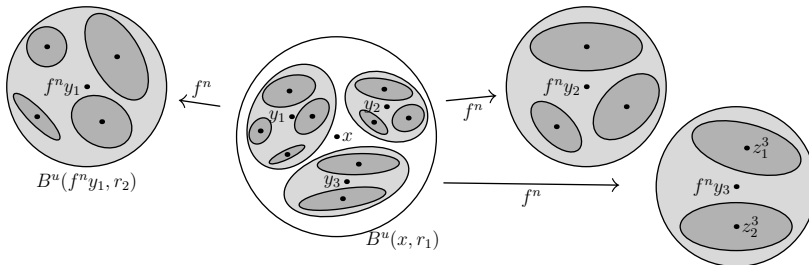


FIGURE 6.1. Proving that $Z_{n+m}^{\text{sep}} \geq C^{-1} Z_n^{\text{sep}} Z_m^{\text{sep}}$

associate a maximal (m, r_2) -separated set $E_m^i = \{z_1^i, \dots, z_{b_i}^i\} \subset B_\Lambda^u(f^n y_i, r_2)$, then pulling back all the points z_j^i gives an $(m + n, r_2)$ -separated set

$$E_{m+n} = \bigcup_{i=1}^a f^{-n}(E_m^i) = \{f^{-n} z_j^i : 1 \leq i \leq a, 1 \leq j \leq b_i\} \subset B_\Lambda^u(x, r_1),$$

so we expect to get an estimate along the lines of

$$\begin{aligned} Z_{m+n}^{\text{sep}} &\geq \sum_{p \in E_{m+n}} e^{S_{m+n}\varphi(p)} = \sum_{p \in E_{m+n}} e^{S_n\varphi(p)} e^{S_m\varphi(f^n(p))} \\ (6.3) \quad &\gtrsim \sum_{i=1}^a \sum_{j=1}^{b_i} e^{S_n\varphi(y_i)} e^{S_m\varphi(z_j^i)} = \sum_{i=1}^a e^{S_n\varphi(y_i)} \left(\sum_{j=1}^{b_i} e^{S_m\varphi(z_j^i)} \right) \\ &\approx \left(\sum_{i=1}^a e^{S_n\varphi(y_i)} \right) Z_m^{\text{sep}} \approx Z_m^{\text{sep}} Z_n^{\text{sep}}, \end{aligned}$$

where we continue to be deliberately vague about the arguments of Z^{sep} . If we can make this rigorous, then a similar argument with spanning sets instead of separated sets will lead to $Z_{m+n}^{\text{span}} \leq C Z_m^{\text{span}} Z_n^{\text{span}}$, which will prove Proposition 6.1.

But how do we make (6.3) rigorous? There are two sources of error which are hinted at by the “ \approx ” symbols.

- (1) Given $p \in E_{m+n}$ and the corresponding $y \in E_n$ (where $f^n(p) \in B_\Lambda^u(f^n y, r_2)$), the approximation on the first line of (6.3) requires us to compare the ergodic sums $S_n\varphi(p)$ and $S_n\varphi(y)$. In particular, we must find a constant Q_u (independent of y, p, n) such that $|S_n\varphi(p) - S_n\varphi(y)| \leq Q_u$ whenever $p \in B_n^u(y, r_2)$.
- (2) The omission of the arguments for Z_n^{sep} obscures the fact that Z_{m+n}^{sep} and Z_m^{sep} in (6.3) both refer to (n, r_2) -separated subsets of $B_\Lambda^u(x, r_1)$, while Z_n^{sep} refers to (n, r_2) -separated subsets of $B_\Lambda^u(f^n y_i, r_2)$. Thus we must control how $Z_n^{\text{sep}}(B_\Lambda^u(x, r_1), \varphi, r_2)$ changes when we fix n and let x, r_1, r_2 vary; in particular, we must find for each r_1, r'_1, r_2, r'_2 a constant C such that for every n, x, y , we have

$$Z_n^{\text{sep}}(B_\Lambda^u(x, r_1), \varphi, r_2) = C^{\pm 1} Z_n^{\text{sep}}(B_\Lambda^u(y, r'_1), \varphi, r'_2).$$

The first source of error described above can be controlled by establishing a generalized version of property (4.4).

Definition 6.5. We say that a potential $\varphi: \Lambda \rightarrow \mathbb{R}$ has the *u-Bowen property* if there is $Q_u > 0$ such that for every $x \in \Lambda$, $n \geq 0$, and $y \in B_n^u(x, \tau) \cap \Lambda$, we have $|S_n\varphi(x) - S_n\varphi(y)| \leq Q_u$. We also say that φ has the *s-Bowen property* if there is $Q_s > 0$ such that for every $x \in \Lambda$, $n \geq 0$, and $y \in B_n^s(x, \tau) = B^s(x, \tau) \cap \Lambda$, we have $|S_n\varphi(x) - S_n\varphi(y)| \leq Q_s$.²⁸

Lemma 6.6. *If $\varphi: \Lambda \rightarrow \mathbb{R}$ is Hölder continuous, then φ has the u-Bowen property and the s-Bowen property.*²⁹

²⁸The asymmetry in the definition comes because $S_n\varphi$ is a forward Birkhoff sum and $B_n^u(x, \tau)$ is defined in terms of forward iterates; one could equivalently define the s-Bowen property in terms of backward Birkhoff sums and s-Bowen balls. The s-Bowen property is needed to control the second source of error described above.

²⁹This is the only place where Hölder continuity is used; in particular, Hölder continuity could be replaced by the u- and s-Bowen properties in all our main results.

Proof. We prove the u -Bowen property; the proof of the s -Bowen property is similar. Given $y \in B_n^u(x, \tau)$, for every $0 \leq k < n$, we have $d(x, y) \leq \tau \lambda^{n-k}$ where $0 < \lambda < 1$ is as in Proposition 3.1(4), so writing β for the Hölder exponent of φ , we have

$$\begin{aligned} |S_n \varphi(x) - S_n \varphi(y)| &\leq \sum_{k=0}^{n-1} |\varphi(f^k x) - \varphi(f^k y)| \leq \sum_{k=0}^{n-1} |\varphi|_\beta d(f^k x, f^k y)^\beta \\ &\leq |\varphi|_\beta \tau^\beta \sum_{k=0}^{n-1} \lambda^{\beta(n-k)} < |\varphi|_\beta \tau^\beta (1 - \lambda^\beta)^{-1} =: Q_u. \quad \square \end{aligned}$$

To control the second source of error described above, the first main idea is that topological transitivity guarantees that for every $\delta > 0$, the images $f^k(B_\Lambda^u(y, \delta))$ eventually come within δ of x , and that the k for which this occurs admits an upper bound that depends only on δ . Then given a spanning set $E \subset B_\Lambda^u(y, \delta)$, the part of the image $f^k(E)$ that lies near x can be moved by holonomy along stable manifolds to give a spanning set in the unstable leaf of x . This is made precise in [CPZ18, Lemma 6.4]. One can use similar arguments to change the scales r_1, r_2 ; for example, if x, y are on the same local unstable leaf and have orbits that remain within r_2 of each other until time n , then they remain within $r_2 \lambda^k$ of each other until time $n - k$; see [CPZ18, §6] for full details.

6.1.3. *Proving Theorem 4.7.* Fix $x \in \Lambda$ and set $X := V_{\text{loc}}^u(x) \cap \Lambda$. We showed in §5.2.3 that m_x^C defines a metric outer measure on X , and hence gives a Borel measure. Note that the final claim in Theorem 4.7 about agreement on intersections is immediate from the definition. Thus it remains to prove that $m_x^C(X) \in [K^{-1}, K]$, where K is independent of x ; this will complete the proof of Theorem 4.7, and will also prove Proposition 5.4. For full details, see [CPZ18, §6.5].

The idea is that it suffices to prove that for a fixed $r > 0$, we have $m_x^C(B_\Lambda^u(x, r))$ uniformly bounded away from 0 and ∞ , since each $V_{\text{loc}}^u(x)$ can be covered with a uniformly finite number of balls $B_\Lambda^u(y, r)$. The upper bound is easier to prove since it only requires that we exhibit a cover satisfying the desired inequality. This is provided by Proposition 6.1, which guarantees existence of an (n, r) -spanning set $E_n \subset B_\Lambda^u(x, r)$ such that

$$\sum_{y \in E_n} e^{S_n \varphi(y)} \leq C e^{nP(\varphi)},$$

and thus (4.13) gives

$$m_x^C(B_\Lambda^u(x, r)) \leq \lim_{n \rightarrow \infty} \sum_{y \in E_n} e^{-nP(\varphi)} e^{S_n \varphi(y)} \leq C.$$

The lower bound is a little trickier since we must obtain a lower bound for an *arbitrary* cover by u -Bowen balls as in (4.13), which are allowed to be of different orders, so we do not immediately get an (n, r) -spanning set for some particular n . This can be resolved by observing that any open cover of $\overline{B_\Lambda^u(x, r)}$ has a finite subcover, so to bound $m_x^C(\overline{B_\Lambda^u(x, r)})$ it suffices to consider covers of the form $\{B_{n_i}^u(y_i, r) : 1 \leq i \leq a, n_i \geq N\}$. Given such a cover, one can take $n = \max(n_1, \dots, n_a)$ and use arguments similar to those in the proof of Proposition 6.1 to cover each $B_{n_i}^u(y_i, r)$ by a union of u -Bowen balls $B_n^u(z_i^j, r)$ ($1 \leq j \leq b_i$)

satisfying

$$\sum_j e^{S_n \varphi(z_i^j)} \leq C' e^{(n-n_i)P(\varphi)} e^{S_{n_i} \varphi(y_i)}$$

for some constant C' that is independent of our choice of covers. Then the set $E = \{z_i^j : 1 \leq i \leq a, 1 \leq j \leq b_i\}$ is (n, r) -spanning for $B_\lambda^u(x, r)$ and satisfies

$$C^{-1} e^{nP(\varphi)} \leq \sum_{z \in E} e^{S_n \varphi(z)} \leq \sum_{i=1}^a C' e^{(n-n_i)P(\varphi)} e^{S_{n_i} \varphi(y_i)}.$$

Dividing through by $e^{nP(\varphi)}$, taking an infimum over all covers, and sending $N \rightarrow \infty$ gives $C^{-1} \leq C' m_x^C(\overline{B_\lambda^u(x, r)})$. Again, full details are in [CPZ18, §6.5].

6.2. Behavior of reference measures under iteration and holonomy.

6.2.1. *Iteration and the u -Gibbs property.* The simplest case of Theorem 4.8 occurs when $\varphi = 0$, so the claim is that $m_{f(x)}^C = e^{P(0)} m_x^C \circ f^{-1}$, which is exactly the scaling property satisfied by the Margulis measures on unstable leaves. Given $E \subset V_{\text{loc}}^u(f(x))$, we see from the relationship $f^{-1} B_n^u(y, r) = B_{n+1}^u(f^{-1}y, r)$ that any cover $\{B_{n_i}^u(y_i, r)\}$ of E leads immediately to a cover $\{B_{n_i+1}^u(f^{-1}y_i, r)\}$ of $f^{-1}E$, and vice versa. Using this bijection in the definition of the reference measures in (4.13), we get

$$\begin{aligned} m_{f(x)}^C(E) &= \lim_{N \rightarrow \infty} \inf \sum_i e^{-n_i P(0)} \\ &= e^{P(0)} \lim_{N \rightarrow \infty} \inf \sum_i e^{-(n_i+1)P(0)} = e^{P(0)} m_x^C(f^{-1}E). \end{aligned}$$

For nonzero potentials one must account for the factor of $e^{S_{n_i} \varphi(x_i)}$ in (4.13). This can be done by partitioning E into subsets E_1, \dots, E_T on which φ is nearly constant, and repeating the above argument on each E_i to get an approximate result that improves to the desired result as $T \rightarrow \infty$; see [CPZ18, §7.1] for details.

Once (1.2) has been proved, we can iterate it to obtain

$$(6.4) \quad m_{f^n(x)}^C(A) = \int_{f^{-n}(A)} e^{nP(\varphi) - S_n \varphi(y)} dm_x^C(y)$$

for all $A \subset V_{\text{loc}}^u(f^n(x))$. Applying this to $A = B^u(f^n(x), \delta) = f^n(B_n^u(x, \delta))$ and using Theorem 4.7 gives a constant $Q_4 = Q_4(\delta)$ such that

$$\begin{aligned} m_x^C(B_n^u(x, \delta)) e^{nP(\varphi) - S_n \varphi(x)} &= e^{\pm Q_u} \int_{B_n^u(x, \delta)} e^{nP(\varphi) - S_n \varphi(y)} dm_x^C(y) \\ &= e^{\pm Q_u} m_{f^n(x)}^C(B^u(f^n(x), \delta)) = e^{\pm Q_u} Q_4^{\pm 1}, \end{aligned}$$

for every x, n , where the first estimate uses the u -Bowen property from Lemma 6.6. This establishes the u -Gibbs property for m_x^C with $Q_1 = Q_4 e^{Q_u}$ and proves Corollary 4.9.

6.2.2. *Holonomy maps.* Given nearby points y, z and sets $E_y \subset V_{\text{loc}}^u(y)$, $E_z \subset V_{\text{loc}}^u(z)$ such that $\pi_{yz}(E_y) = E_z$ (with respect to some rectangle), we observe that every cover of E_y by u -Bowen balls $\{B_{n_i}^u(x_i, r)\}$ produces a cover of E_z by the images $\{\pi_{yz} B_{n_i}^u(x_i, r)\}$. If y, z are close enough to each other to guarantee that

$$(6.5) \quad \pi_{yz} B_{n_i}^u(x_i, r) \subset B_{n_i}^u(x_i, 2r)$$

for each i , then we get $E_z \subset \bigcup_i B_{n_i}^u(x_i, 2r)$. Fixing $k \in \mathbb{N}$ such that each $x \in \Lambda$ has $B_\Lambda^u(x, 2r) \subset \bigcup_{j=1}^k B_\Lambda^u(x^j, r)$ for some x^1, \dots, x^k , we see that $E_z \subset \bigcup_{i,j} B_{n_i}^u(x_i^j, r)$, and thus (4.13) gives

$$m_z^C(E_z) \leq \sum_{i,j} e^{-n_i P(\varphi)} e^{S_{n_i} \varphi(x_i^j)} \leq k \sum_i e^{-n_i P(\varphi)} e^{S_{n_i} \varphi(x_i) + Q_u}.$$

Taking an infimum and then a limit gives $m_z^C(E_z) \leq k e^{Q_u} m_y^C(E_y)$.

In general, if y, z lie close enough for holonomy maps to be defined but not close enough for (6.5) to hold, then we can iterate E_y, E_z forward until some time n at which $f^n y, f^n z$ are close enough for the previous part to work, and use Theorem 4.8 to get (assuming without loss of generality that $E_y \subset B_n^u(y, \tau)$, and similarly for E_z)

$$\begin{aligned} m_z^C(E_z) &= \int_{f^n(E_z)} e^{-nP(\varphi) + S_n \varphi(f^{-n}x)} dm_{f^n z}^C(x) \\ &\leq k e^{Q_u} \int_{f^n(E_y)} e^{-nP(\varphi) + S_n \varphi(f^{-n}x') + Q_s} dm_{f^n y}^C(x') = k e^{Q_u + Q_s} m_y^C(E_y), \end{aligned}$$

where the inequality uses the result from the previous paragraph. Since the roles of y, z were symmetric, this proves Theorem 4.10 with $Q_2 = k e^{Q_u + Q_s}$. See [CPZ18, §7.3] for a more detailed version of this argument.

6.3. Geometric construction of equilibrium states. Now that we have established the basic properties of the reference measures m_x^C associated to a Hölder continuous potential function φ , the steps in the geometric construction of the unique equilibrium state μ_φ are as follows.

- (1) Prove that every weak* limit point μ of the sequence of probability measures $\mu_n = \frac{1}{n} \sum_{k=0}^{n-1} f_*^k m_x^C / m_x^C(V_{\text{loc}}^u(x))$ is an invariant measure whose conditional measures satisfy part (3) of Theorem 4.11; in particular, they are equivalent to the reference measures m_y^C .
- (2) Use this to deduce that any such μ satisfies part (2) of Theorem 4.11, namely:
 - (a) the conditional measures of μ are absolutely continuous with respect to stable holonomies, and therefore μ is ergodic by the Hopf argument (Proposition 4.5);
 - (b) μ gives positive weight to every open set in Λ ;
 - (c) the u -Gibbs property of the reference measures implies the Gibbs property (3.15) for μ ; and
 - (d) μ is the unique equilibrium state for φ by Proposition 3.11.
- (3) Observe that each μ_n is a Borel probability measure on Λ , and thus every subsequence has a subsubsequence that converges in the weak* topology to a Borel probability measure μ , which must be the unique equilibrium state μ_φ by the previous step. Since every subsequence of μ_n has a subsubsequence converging to μ_φ , it follows that the sequence itself converges to this limit, which establishes part (1) of Theorem 4.11.

The first step takes most of the work; once it is done, parts (a)–(c) of the second step only require short arguments that leverage the properties already established, part (d) of the second step merely consists of observing that μ satisfies the hypotheses of Proposition 3.11, and the third step is completely contained in the description

in (3). Thus we outline here the argument for the first step and parts (a)–(c) of the second step, referring once more to [CPZ18] for complete details.

6.3.1. *Conditional measures of limiting measures.* In order to understand the conditional measures of $\mu = \lim_{k \rightarrow \infty} \mu_{n_k}$, we start by studying the conditional measures of μ_n . Given $x \in \Lambda$ and $n \in \mathbb{N}$, the iterate $f_*^n m_x^C$ is supported on $W_n = f^n(V_{\text{loc}}^u(x) \cap \Lambda)$, and given any $y \in W_n$, we can iterate the formula from Theorem 4.8 and obtain

$$(6.6) \quad \frac{d((f_*^n m_x^C)|_{V_{\text{loc}}^u(y)})}{dm_y^C}(z) = e^{-nP(\varphi) + S_n \varphi(f^{-n}z)} =: g_n(z)$$

for every $z \in W_n \cap V_{\text{loc}}^u(y)$. One can show that $g_n \rightarrow 0$ as $n \rightarrow \infty$, so it is convenient to write $\rho_n^y(z) := g_n(z)/g_n(y)$, and Lemma 6.6 gives

$$(6.7) \quad \rho_n^y(z) = e^{S_n \varphi(f^{-n}z) - S_n \varphi(f^{-n}y)} \in [e^{-Q_u}, e^{Q_u}].$$

These functions describe the conditional measures of $f_*^n m_x^C$. Indeed, given a rectangle R , choose $y_1, \dots, y_a \in f^n(V_{\text{loc}}^u(x)) \cap R$ such that $f^n(V_{\text{loc}}^u(x)) \cap R \subset \bigcup_{i=1}^a V_{\text{loc}}^u(y_i)$, as in Figure 6.2. Then for every Borel set $E \subset R$, we have³⁰

$$(6.8) \quad f_*^n m_x^C(E) = \sum_{i=1}^a \int_E g_n(z) dm_{y_i}^C(z) = \sum_{i=1}^a g_n(y_i) \int_E \rho_n^{y_i}(z) dm_{y_i}^C(z).$$

In other words, one can write $f_*^n m_x^C|_R$ as a linear combination of the measures $\rho_n^{y_i} dm_{y_i}^C$ associated to the *standard pairs*³¹ $(V_{\text{loc}}^u(y_i), \rho_n^{y_i})$, with coefficients given by $g_n(y_i)$. This immediately implies that the conditional measures of μ_n on local unstable leaves are absolutely continuous with respect to the reference measures m_x^C , with densities bounded away from 0 and ∞ .

To go further, we need the following characterization of the conditional measures, which is an immediate consequence of [EW11, Corollary 5.21].³²

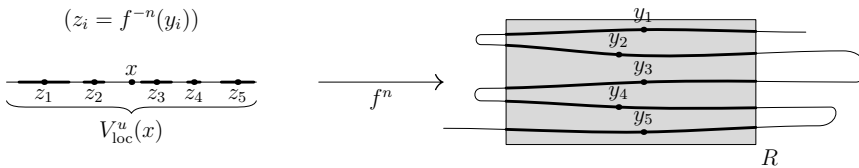
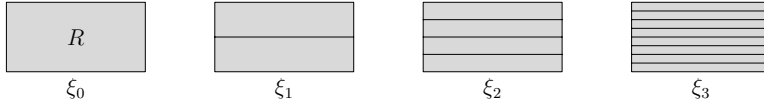


FIGURE 6.2. Studying $f_*^n m_x^C$ on a rectangle R

³⁰There is a small technical issue here, namely that there may be some y_i at which W_n does not cross R completely, and so the integral in (6.8) actually gives too large a value. However, this can only occur if $z_i = f^{-n}(y_i)$ is very close to the boundary of $V_{\text{loc}}^u(x)$, and the contribution of such points is negligible in the limit; see [CPZ18].

³¹Standard pairs consisting of a local leaf $V_{\text{loc}}^u(y)$ and a density function ρ were introduced by Chernov and Dolgopyat in [CD09] to study stochastic properties of dynamical systems. They are also used in constructing SRB measures for some dynamical systems with weak hyperbolicity; see [CDP16].

³²See [PS82] for the analogous argument controlling the conditionals of μ when φ is the geometric potential and the reference measure is leaf volume. In that setting, the role of Proposition 6.7 here is played by [PS82, Lemma 13].

FIGURE 6.3. A refining sequence of partitions of R

Proposition 6.7. *Let μ be a finite Borel measure on Λ , and let $R \subset \Lambda$ be a rectangle with $\mu(R) > 0$. Let $\{\xi_\ell\}_{\ell \in \mathbb{N}}$ be a refining sequence of finite partitions of R that converge to the partition ξ into local unstable sets $V_R^u(y) = V_{\text{loc}}^u(y) \cap R$. Then there is a set $R' \subset R$ with $\mu(R') = \mu(R)$ such that for every $y \in R'$ and every continuous $\psi: R \rightarrow \mathbb{R}$, we have*

$$(6.9) \quad \int_{V_R^u(y)} \psi(z) d\mu_y^u(z) = \lim_{\ell \rightarrow \infty} \frac{1}{\mu(\xi_\ell(y))} \int_{\xi_\ell(y)} \psi(z) d\mu(z),$$

where $\xi_n(y)$ denotes the element of the partition ξ_ℓ that contains y .

Now the proof of part (3) of Theorem 4.11 goes as follows. Given a rectangle $R \subset \Lambda$ with $\mu(R) > 0$, let ξ_ℓ be a refining sequence of finite partitions of R such that for every $y \in R$ and $\ell \in \mathbb{N}$, the set $\xi_\ell(y)$ is a rectangle, and $\bigcap_{\ell \in \mathbb{N}} \xi_\ell(y) = V_R^u(y)$, as in Figure 6.3. Let $R' \subset R$ be the set given by Proposition 6.7. We prove that (4.14) holds for each $y \in R'$ by showing that for every positive continuous function $\psi: R \rightarrow \mathbb{R}$, we have

$$(6.10) \quad \int_{V_R^u(y)} \psi d\mu_y^u = \frac{Q_3^{\pm 1}}{m_y^C(V_R^u(y))} \int_{V_R^u(y)} \psi dm_y^C,$$

where Q_3 is a constant that is independent of ψ . To this end, we need to compare $\mu(\xi_\ell(y))$ and $\int_{\xi_\ell(y)} \psi(z) d\mu(z)$ and then apply (6.9). We see from (6.7) and (6.8) that for each $j \in \mathbb{N}$, there is a finite set $Y_j \subset R$ such that

$$(6.11) \quad \int_{\xi_\ell(y)} \psi d(f_*^n m_x^C) = \sum_{p \in Y_j} g_j(p) \int_{V_R^u(p)} \psi(z) e^{\pm Q_u} dm_p^C(z).$$

Given $p \in \xi_\ell(y)$, Theorem 4.10 gives

$$\int_{V_R^u(p)} \psi dm_p^C = Q_2^{\pm 1} \int_{V_R^u(y)} \psi(\pi_{py} z') dm_y^C(z') = (2Q_2)^{\pm 1} \int_{V_R^u(y)} \psi dm_y^C$$

whenever p, y are sufficiently close that $\psi(\pi_{py} z') = 2^{\pm 1} \psi(z')$ for all $z' \in V_R^u(y)$. Thus for all sufficiently large ℓ , (6.11) gives

$$\int_{\xi_\ell(y)} \psi d(f_*^n m_x^C) = e^{\pm Q_u} (2Q_2)^{\pm 1} \left(\sum_{p \in Y_j} g_j(p) \right) \int_{V_R^u(y)} \psi dm_y^C.$$

Averaging over $0 \leq j < n_k$ and sending $k \rightarrow \infty$ gives

$$(6.12) \quad \int_{\xi_\ell(y)} \psi d\mu = (2Q_2 e^{Q_u})^{\pm 1} \left(\overline{\lim}_{k \rightarrow \infty} \frac{1}{n_k} \sum_{j=0}^{n_k-1} \sum_{p \in Y_j} g_k(p) \right) \int_{V_R^u(y)} \psi dm_y^C.$$

When $\psi \equiv 1$, this gives

$$(6.13) \quad \mu(\xi_\ell(y)) = (2Q_2 e^{Q_u})^{\pm 1} \left(\overline{\lim}_{k \rightarrow \infty} \frac{1}{n_k} \sum_{j=0}^{n_k-1} \sum_{p \in Y_j} g_k(p) \right) m_y^C(V_R^u(y)).$$

Dividing (6.12) by (6.13), sending $\ell \rightarrow \infty$, and using (6.9) yields

$$\int_{V_R^u(y)} \psi d\mu_y^u = (2Q_2 e^{Q_2 u})^{\pm 2} \frac{1}{m_y^C(V_R^u(y))} \int_{V_R^u(y)} \psi dm_y^C.$$

Since $\psi > 0$ was arbitrary, this proves (4.14), modulo some minor technical issues around the boundary of $V_{\text{loc}}^u(x)$ that are dealt with in [CPZ18].

6.3.2. Other properties of limiting measures. Throughout this section, μ will denote an arbitrary f -invariant Borel probability measure on Λ that satisfies (4.14), so that the conditional measures μ_y^u are equivalent to the reference measures m_y^C . By the previous section, this includes every limit point of the sequence μ_n .

We first observe that by Theorem 4.10 and (4.14), for every rectangle R with $\mu(R) > 0$, μ -a.e. $y, z \in R$, and every $A \subset V_R^u(z)$, we have

$$(6.14) \quad \mu_y^u(\pi_{zy}A) = Q_3^{\pm 1} m_y^C(\pi_{zy}A) / m_y^C(R) = Q_3^{\pm 1} Q_2^{\pm 2} m_z^C(A) / m_z^C(R) = (Q_3 Q_2)^{\pm 2} \mu_z^u(A).$$

In particular, holonomy maps along stable manifolds are absolutely continuous with respect to the conditional measures μ_y^u , and thus by the standard Hopf argument (Proposition 4.5), μ is ergodic.

Now we prove that μ is fully supported and satisfies the Gibbs property. Let $\delta > 0$ be small enough that for every $x \in \Lambda$, the rectangle

$$R(x, \delta) := [\overline{B_\Lambda^u(x, \delta)}, \overline{B_\Lambda^s(x, \delta)}] = \{[y, z] : y \in \overline{B_\Lambda^u(x, \delta)}, z \in \overline{B_\Lambda^s(x, \delta)}\}$$

is well defined, as in (3.2). Given $n \in \mathbb{N}$, consider the rectangle

$$R_n(x, \delta) := [\overline{B_n^u(x, \delta)} \cap \Lambda, \overline{B_n^s(x, \delta)}] \subset R(x, \delta).$$

It is shown in [CPZ18, Lemma 8.3] that for every $\delta > 0$, there are $\delta_1, \delta_2 > 0$ such that

$$(6.15) \quad R_n(x, \delta_1) \subset B_n(x, \delta) \subset R_n(x, \delta_2)$$

for every $x \in \Lambda$ and $n \in \mathbb{N}$; thus to prove the Gibbs property (3.14) it suffices to establish the corresponding bounds on $\mu(R_n(x, \delta))$.

Lemma 6.8. *Given $\delta > 0$, there is $Q_5 > 0$ such that for every x, δ, n as above, we have*

$$(6.16) \quad \mu(R_n(x, \delta)) = Q_5^{\pm 1} e^{-nP(\varphi) + S_n \varphi(x)} \mu(R(x, \delta)).$$

Proof. Writing μ_y^u for the conditional measures of μ on unstable leaves in $R(z, \delta)$, we have

$$\begin{aligned} \mu(R_n(x, \delta)) &= \int_{R(x, \delta)} \mu_y^u(R_n(x, \delta)) d\mu(y) \\ &= Q_3^{\pm 1} \int_{R(x, \delta)} \frac{m_y^C(R_n(x, \delta))}{m_y^C(R(x, \delta))} d\mu(y) \\ &= (KQ_3)^{\pm 1} \int_{R(x, \delta)} m_y^C(\pi_{xy} \overline{B_n^u(x, \delta)}) d\mu(y) \\ &= (KQ_3 Q_2)^{\pm 1} m_x^C(\overline{B_n^u(x, \delta)}) \mu(R(x, \delta)), \end{aligned}$$

where the first equality uses the definition of conditional measures, the second uses (4.14), the third uses Theorem 4.7, and the fourth uses Theorem 4.10. Since

$B_n^u(x, \delta) \subset \overline{B_n^u(x, \delta)} \subset B_n^u(x, 2\delta)$, the result follows from the u -Gibbs property of m_x^c . \square

Lemma 6.9 ([CPZ18, Lemma 8.4]). *For every sufficiently small $\delta > 0$, there is $\delta' > 0$ such that for every $z \in \Lambda$ and $x \in R(z, \delta')$, we have $R(z, \delta') \subset R(x, \delta)$.*

Lemma 6.10. *If $y \in \Lambda$ has a backward orbit that is dense in Λ , then $\mu(R(y, \delta)) > 0$ for all $\delta > 0$.*

Proof. Let $\delta' > 0$ be as in Lemma 6.9. Since Λ is compact, there is a finite set $E \subset \Lambda$ such that $\bigcup_{z \in E} R(z, \delta') = \Lambda$, and thus there is $z \in E$ with $\mu(R(z, \delta')) > 0$. Since the backward orbit of y is dense, there is $n \geq 0$ such that $x := f^{-n}(y) \in R(z, \delta')$. By Lemma 6.9 and our choice of x , we have

$$\mu(R(x, \delta)) \geq \mu(R(z, \delta')) > 0.$$

By Lemma 6.8, we conclude that $\mu(R_n(x, \delta)) > 0$. Moreover, we have

$$f^n R_n(x, \delta) = f^n [B_n^u(x, \delta) \cap \Lambda, B_\Lambda^s(x, \delta)] \subset [B_\Lambda^u(y, \delta), B_\Lambda^s(y, \delta)] = R(y, \delta),$$

where we use the fact that $\|Df|_{E^s}\| \leq 1$. Since μ is f -invariant, this gives $\mu(R(y, \delta)) \geq \mu(R_n(x, \delta)) > 0$. \square

Since f is topologically transitive on Λ , every (relatively) open set in Λ contains a set of the form $R(y, \delta)$ where y has a dense backward orbit. Thus Lemma 6.10 implies that μ is fully supported on Λ .

Finally, we deduce the Gibbs property (3.15) for μ as follows. Given $\delta > 0$, let $\delta' > 0$ be as in Lemma 6.9, and once again let $E \subset \Lambda$ be a finite set with $\bigcup_{z \in E} R(z, \delta') = \Lambda$. Since μ is fully supported, we have $\eta := \min_{z \in E} \mu(R(z, \delta')) > 0$. Now given any $x \in \Lambda$, we have $x \in R(z, \delta')$ for some $z \in E$, and thus Lemma 6.9 gives

$$\mu(R(x, \delta)) \geq \mu(R(z, \delta')) \geq \eta.$$

In particular, $\eta \leq \mu(R(x, \delta)) \leq 1$ for every $x \in \Lambda$, and then the Gibbs property (3.15) follows immediately from Lemma 6.8 and (6.15).

6.3.3. Local product structure. To prove Corollary 4.12, we first observe that (6.14) gives $\pi_{yp}^* \mu_p^u \sim \mu_y^u$ for μ -a.e. $p, y \in R$, which is the first claim. For the second claim, we define a function $h: R \rightarrow (0, \infty)$ by $h(z) = \frac{d\mu_z^u}{d(\pi_{zp}^* \mu_p^u)}(z)$, so that (3.4) gives

$$\begin{aligned} \mu(E) &= \int_{V_R^s(p)} \int_{V_R^u(y)} \mathbf{1}_E(z) d\mu_y^u(z) d\tilde{\mu}_p(y) \\ (6.17) \quad &= \int_{V_R^s(p)} \int_{V_R^u(y)} \mathbf{1}_E(z) h(z) d(\pi_{zp}^* \mu_p^u)(z) d\tilde{\mu}_p(y) \\ &= \int_E h(z) d(\mu_p^u \otimes \tilde{\mu}_p)(z) \end{aligned}$$

for every measurable $E \subset R$. For the third claim, we observe that (6.17) gives

$$\mu(E) = \int_E h(z) d(\mu_p^u \otimes \tilde{\mu}_p)(z) = \int_{V_R^u(p)} \int_{V_R^s(y)} \mathbf{1}_E(z) h(z) d(\pi_{yp}^* \tilde{\mu}_p)(z) d\mu_p^u(y),$$

and since μ_y^s is uniquely determined up to a scalar (for μ -a.e. y) by the condition that

$$\mu(E) = \int_{V_R^u(p)} \int_{V_R^s(y)} \mathbf{1}_E(z) d\mu_y^s(z) d\nu(y)$$

for some measure ν on $V_R^u(p)$, we conclude that $d\mu_y^s = h d(\pi_{yp}^* \tilde{\mu}_p)$.

6.4. **Proof of Theorem 4.1.** We have that for any Borel subset $E \subset Y$

$$(6.18) \quad \mu(E) = \int_{\tilde{Y}} \int_W \mathbf{1}_E(z) d\mu_W^\xi(z) d\tilde{\mu}(W).$$

Without loss of generality we may assume that μ_W^ξ is normalized, so that $\mu_W^\xi(W) = 1$. Consider the set B of all Birkhoff generic points $x \in X$, for which

$$(6.19) \quad \lim_{n \rightarrow \infty} \frac{1}{n} \sum_{k=0}^{n-1} h(f^k(x)) = \int_X h d\mu$$

for every continuous function h on X . Since μ is ergodic, we have that B has full measure in Y . By (6.18), there is a set $D \subset \tilde{Y}$ such that $\tilde{\mu}(\tilde{Y} \setminus D) = 0$, and for every $W \in D$ we have $\mu_W^\xi(W \setminus B) = 0$. Given any such W and any measure $\nu \ll \mu_W^\xi$, we have $\nu(X \setminus B) = 0$. Then Theorem 4.1 is a consequence of the following general result.

Proposition 6.11. *Let X be a compact metric space, let $f: X \rightarrow X$ be a continuous map, and let μ be an f -invariant Borel probability measure on X . Let B be the set of Birkhoff generic points for μ , and let ν be any probability measure on X such that $\nu(B) = 1$. Then the sequence of measures $\nu_n := \frac{1}{n} \sum_{k=0}^{n-1} f_*^k \nu$ converges in the weak* topology to the measure μ .*

Before proving Proposition 6.11, we note that μ is not required to be ergodic. In the case when μ is ergodic, Birkhoff's theorem gives $\mu(B) = 1$, so that in particular B is nonempty. For nonergodic μ , the set B can be either empty or nonempty.

Proof of Proposition 6.11. Let κ be a weak* limit point of the sequence ν_n , so that there is a subsequence $\{n_\ell\}_{\ell \in \mathbb{N}}$ such that for every continuous function h on X , we have

$$(6.20) \quad \int_X h d\kappa = \lim_{\ell \rightarrow \infty} \int_X h d\nu_{n_\ell} = \lim_{\ell \rightarrow \infty} \int \frac{1}{n_\ell} \sum_{k=0}^{n_\ell-1} h \circ f^k d\nu.$$

We show that $\kappa \leq \mu$, which implies that $\kappa = \mu$ since both are probability measures. It suffices to show that $\int h d\kappa \leq \int h d\mu$ for every nonnegative continuous function h .

Fix h as above. Given $N \in \mathbb{N}$ and $\varepsilon > 0$, let

$$B_N(\varepsilon) := \left\{ x \in B : \left| \frac{1}{n} S_n h(x) - \int h d\mu \right| < \varepsilon \text{ for all } n \geq N \right\}.$$

Then for every $\varepsilon > 0$ we have $\bigcup_{N \in \mathbb{N}} B_N(\varepsilon) = B$, hence there is N_ε such that $\nu(B \setminus B_{N_\varepsilon}(\varepsilon)) < \varepsilon$. By (6.20) we can choose $n_\ell > N_\varepsilon$ such that

$$\begin{aligned} \int_X h d\kappa &\leq \varepsilon + \int \frac{1}{n_\ell} S_{n_\ell} h d\nu = \varepsilon + \int_{B_N(\varepsilon)} \frac{1}{n_\ell} S_{n_\ell} h d\nu + \int_{B \setminus B_N(\varepsilon)} \frac{1}{n_\ell} S_{n_\ell} h d\nu \\ &\leq 2\varepsilon + \int h d\mu + \nu(B \setminus B_N(\varepsilon)) \|h\| < \varepsilon(2 + \|h\|) + \int h d\mu. \end{aligned}$$

Since $\varepsilon > 0$ was arbitrary, this completes the proof of the proposition. \square

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